

# Neutron stars and nuclear symmetry energy

Hajime SOTANI (YITP, Kyoto U.)

HS, Nakazato, Iida, Oyamatsu, PRL (2012)

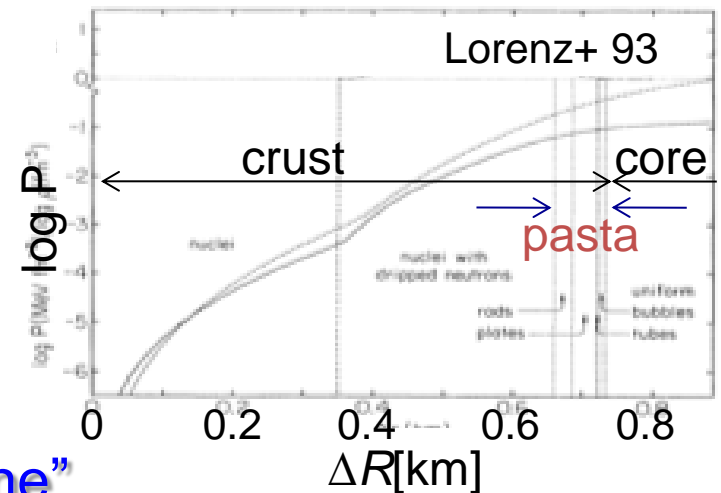
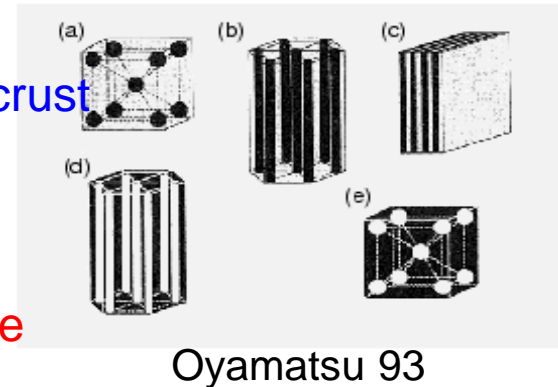
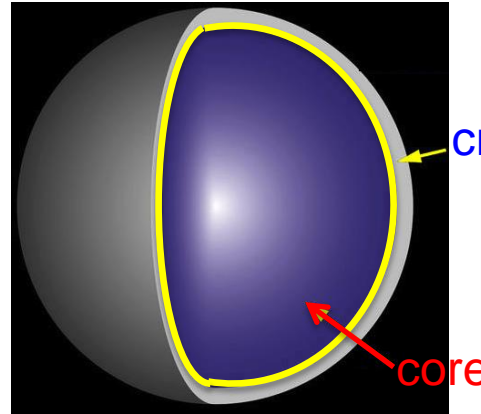
HS, Nakazato, Iida, Oyamatsu, MNRASL (2013)

HS, Nakazato, Iida, Oyamatsu, MNRAS (2013)

# neutron stars

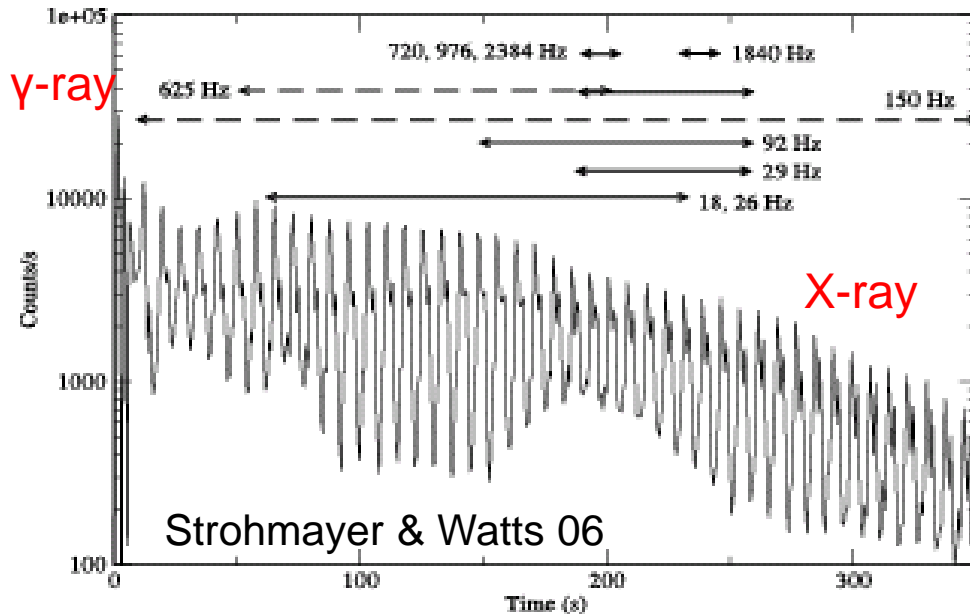
- Structure of NS
  - solid layer (crust)
  - nonuniform structure (pasta)
  - fluid core (uniform matter)
- Crust thickness  $\lesssim 1\text{km}$ 
  - Determination of EOS for high density region could be quite difficult on Earth
- Constraint on EOS via observations of neutron stars
  - stellar mass and radius
  - stellar oscillations (& emitted GWs)

“(GW) asteroseismology”
- NS can be considered as “Rosetta stone” to see physics in ultra-high density region.



# QPOs in SGRs

- Quasi-periodic oscillations (QPOs) in afterglow of giant flares from soft-gamma repeaters (SGRs)
  - SGR 0526-66 (5<sup>th</sup>/3/1979) : 43 Hz
  - SGR 1900+14 (27<sup>th</sup>/8/1998) : 28, 54, 84, 155 Hz
  - SGR 1806-20 (27<sup>th</sup>/12/2004) : 18, 26, 30, 92.5, 150, 626.5, 1837 Hz  
(Barat+ 1983, Israel+ 05, Strohmayer & Watts 05, Watts & Strohmayer 06)



- Crustal torsional oscillation ?
- Magnetic oscillations ?
- Asteroseismology
  - stellar properties  
( $M$ ,  $R$ ,  $B$ , EOS ...)

# torsional oscillations

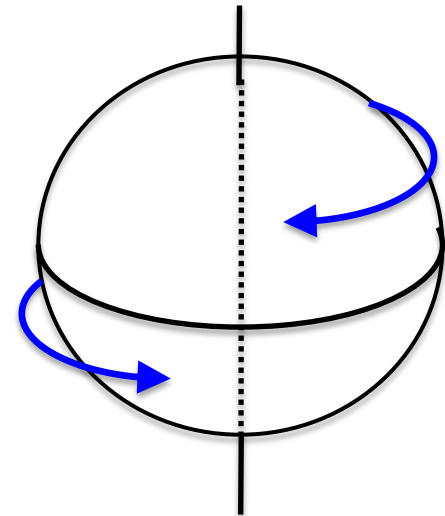
- axial parity oscillations
  - incompressible
  - no density perturbations

- in Newtonian case

(Hansen & Ciolf 1980)

$$\ell t_0 \sim \frac{\sqrt{\ell(\ell+1)\mu/\rho}}{2\pi R} \sim 16\sqrt{\ell(\ell+1)} \text{ Hz} \quad \ell t_n \sim \frac{\sqrt{\mu/\rho}}{2\Delta r} \sim 500 \times n \text{ Hz}$$

- $\mu$ : shear modulus
- frequencies  $\propto$  shear velocity  $v_s = \sqrt{m/r}$
- overtones depend on crust thickness
- effect of magnetic field
  - frequencies become larger  
(Sotani+07, Gabler+13)



# EOS for curst region

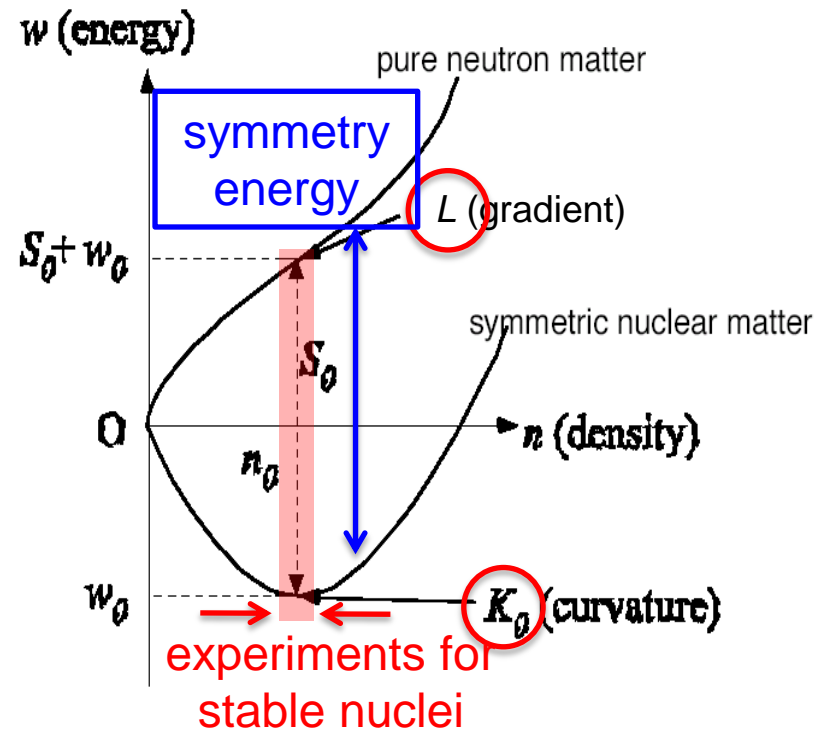
(Oyamatsu & Iida 03, 07)

- Bulk energy per nucleon near the saturation point of symmetric nuclear matter at zero temperature;

$$w = w_0 + \frac{K_0}{18n_0^2}(n - n_0)^2 + \left[ S_0 + \frac{L}{3n_0}(n - n_0) \right] \alpha^2$$

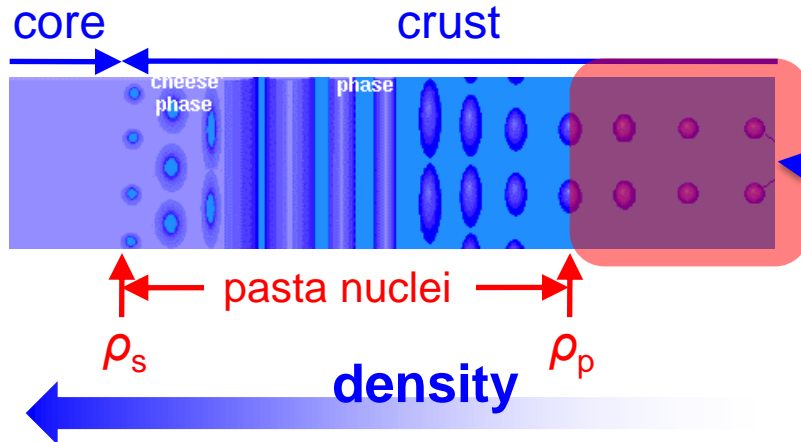
$K_0$  incompressibility
 $L$  symmetry parameter

- Calculations of the optimal density distribution of stable nuclei within Thomas Fermi theory.
- phenomenological, but cover the experimental data for stable nuclei.
- $K_0$  &  $L$  are associated with stiffness EOS of nuclear matter



# what we do

- EOS for core region is still uncertain. (cf. Steiner & Watts 09)
- To prepare the crust region, we integrate from  $r=R$ .
  - $M, R$  : parameters for stellar properties
  - $L, K_0$  : parameters for crust EOS (Oyamatsu & Iida 03, 07)
- In crust region, torsional oscillations are calculated.
  - considering the shear only in spherical nuclei.
  - frequency of fundamental oscillation  $\propto v_s$  ( $v_s^2 \sim \mu/H$ )
- Comparing frequencies with QPOs, we will put a constraint on EOS parameter.



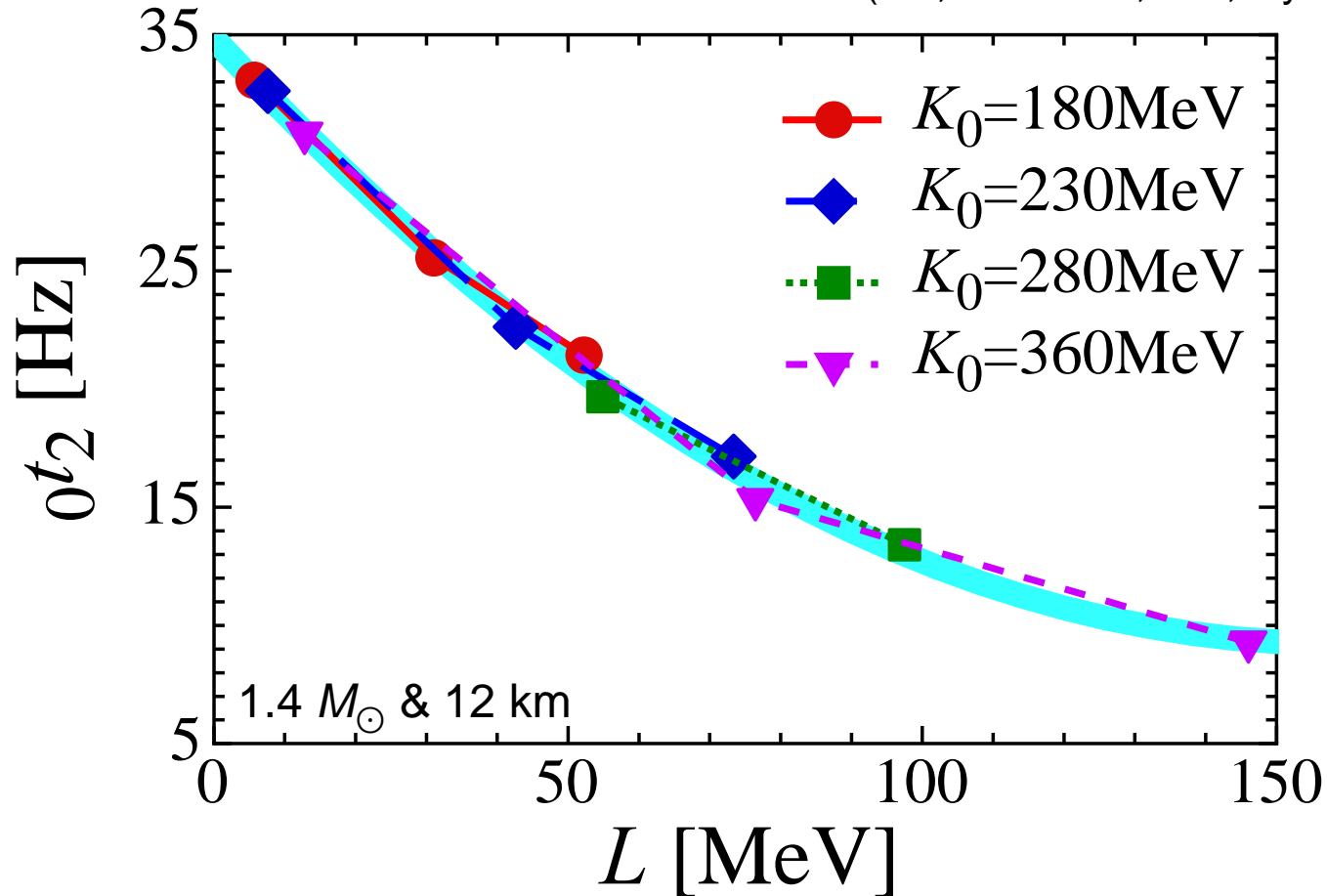
for bcc lattice (Strohmayer+ 91)

$$\mu = 0.1194 \frac{n_i (Ze)^2}{a}$$

$n_i$  : number density of quark droplet  
 $Z$  : charge of quark droplet  
 $a$  : Wigner-Seitz radius

# torsional oscillations

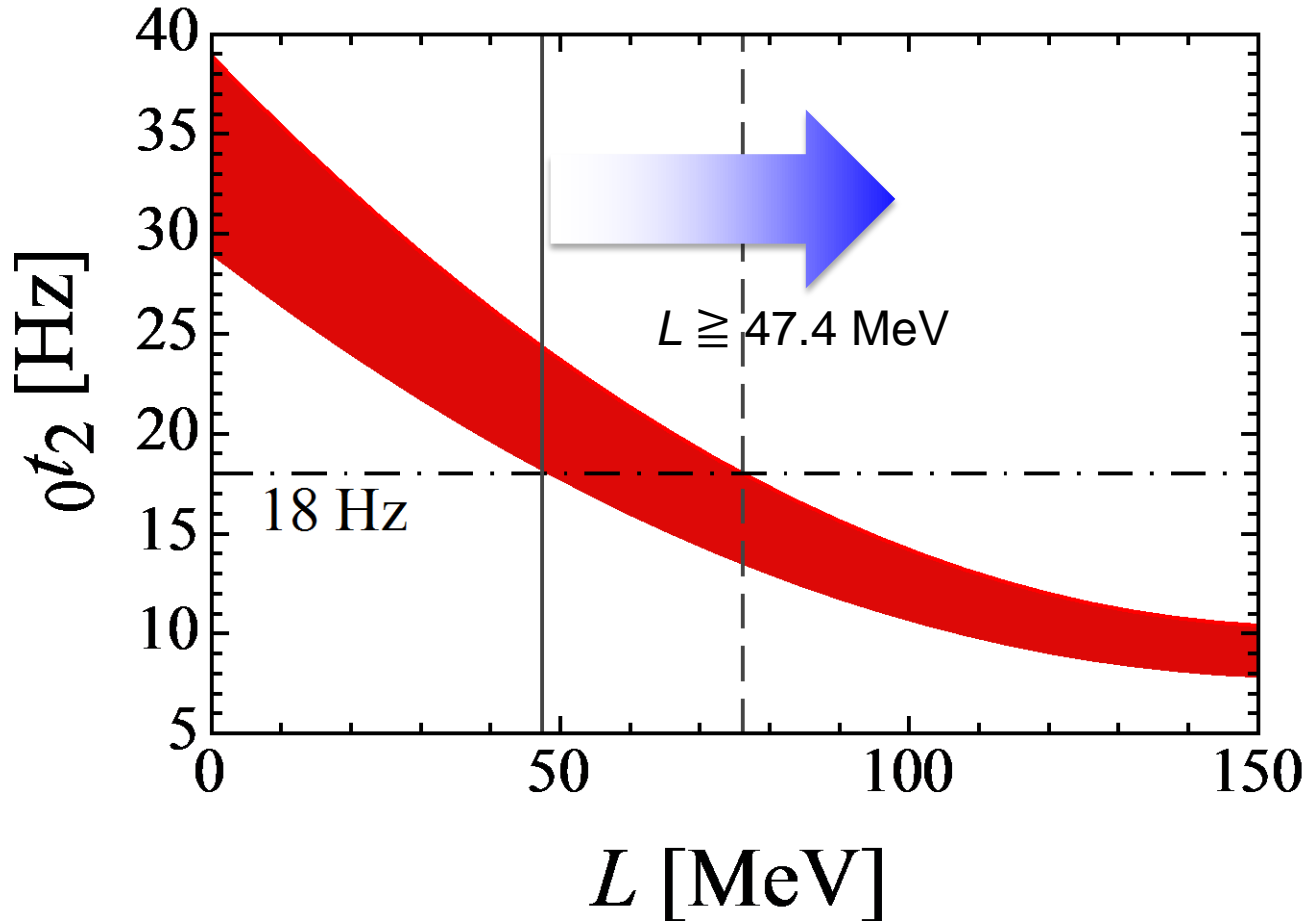
(HS, Nakazato, Iida, Oyamatsu13)



→ almost independent of the incompressibility  $K_0$

# robust constraint on $L$

$10 \text{ km} \leq R \leq 14 \text{ km} \ \& \ 1.4 \leq M/M_{\odot} \leq 1.8$





# effect of superfluidity

- $\rho \gtrsim 4 \times 10^{11} \text{ g/cm}^3$ ; neutrons start to drip out of nuclei
  - some of them play as superfluid
  - how many fraction of dripped neutrons behave as superfluid ?
  - major parts may be locked to the motion of protons in nuclei (Chamel 12)
  - depending on density,  $N_s/N_d \simeq 10 - 30\% @ n_b \sim 0.01 - 0.4 n_0$
- since torsional oscillations are transverse, superfluid neutrons can not contribute to such oscillations.

– one show introduce the effective enthalpy

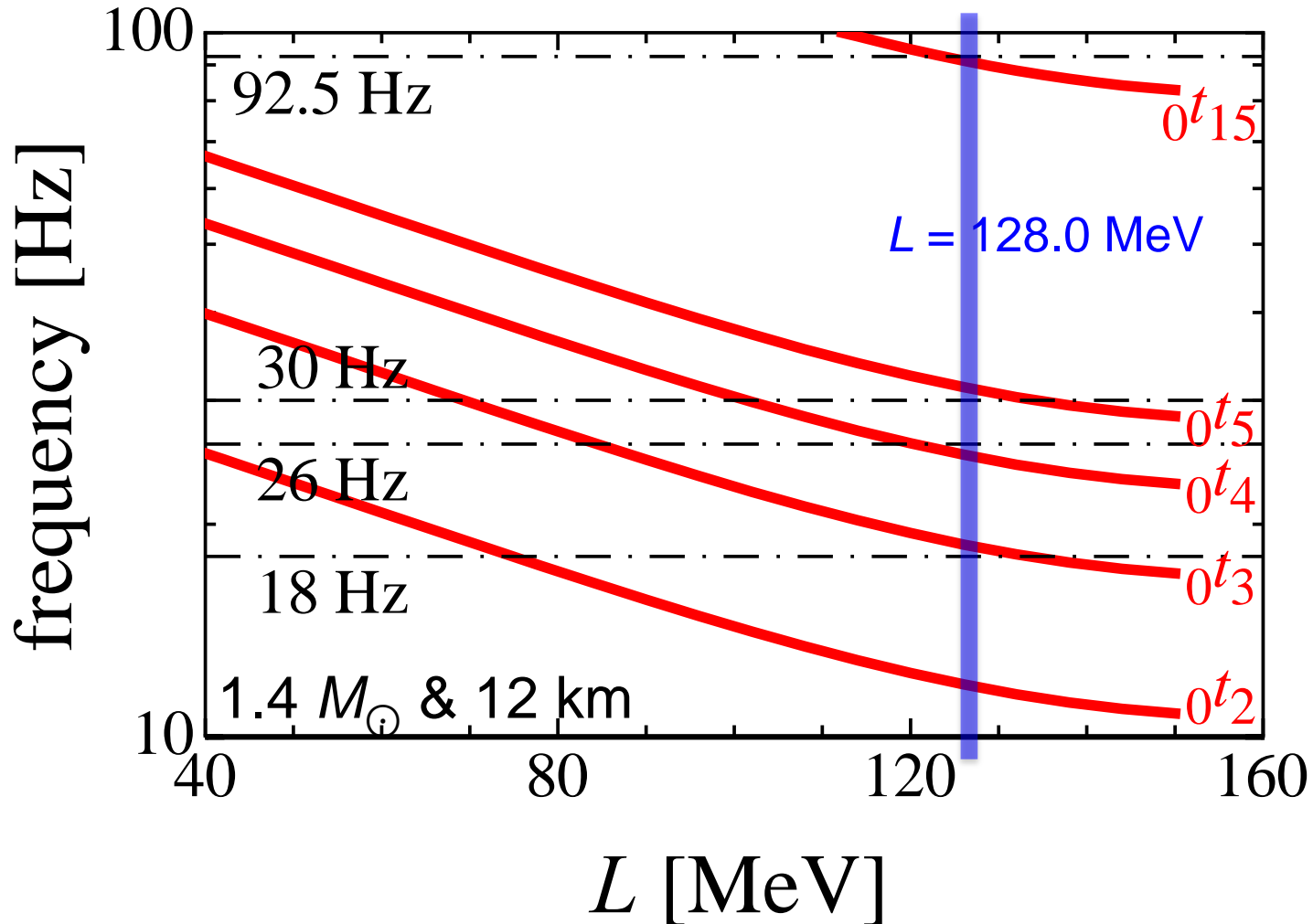
– at zero-temperature,  $\mu_b = H / n_b$



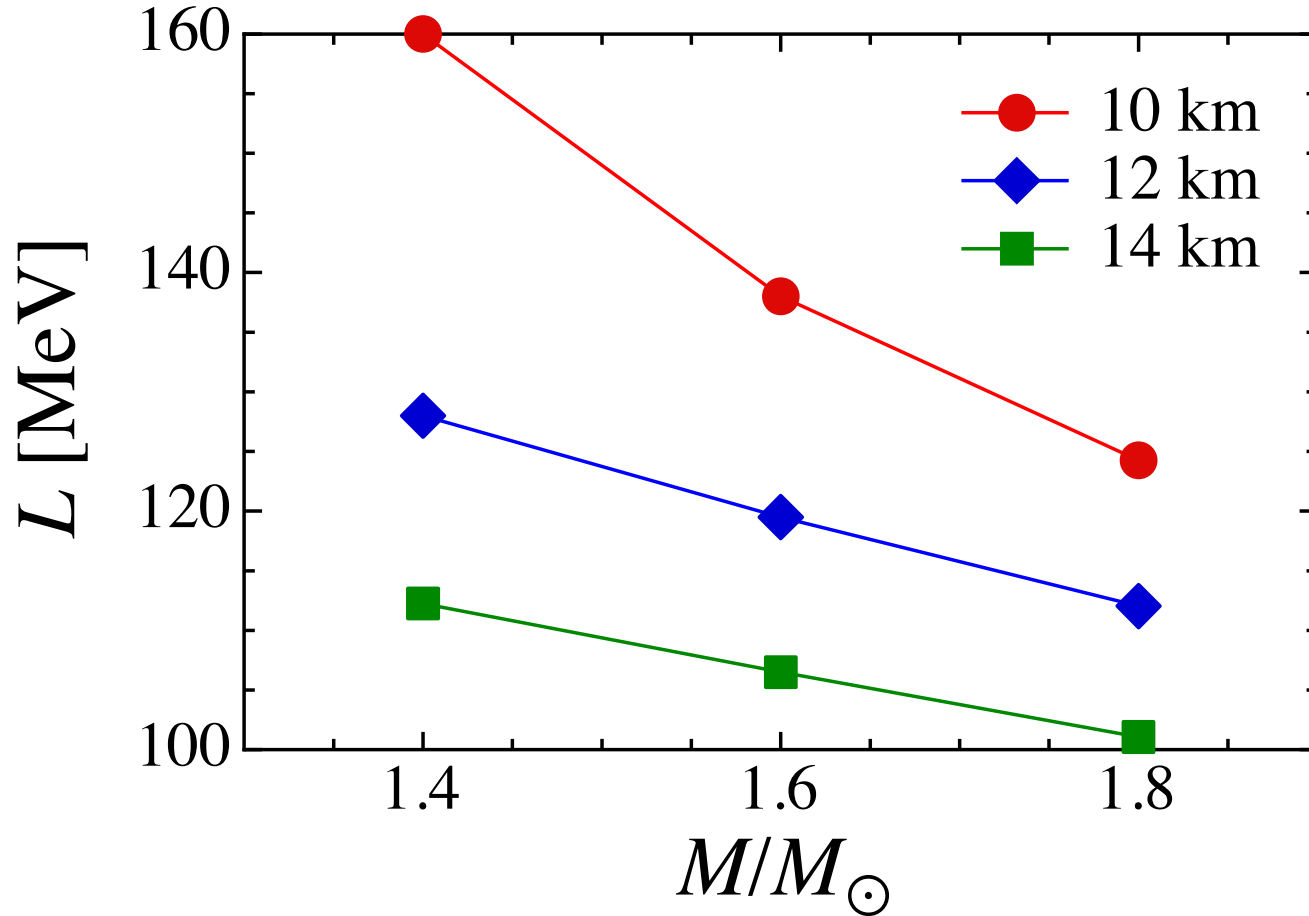
$$\bar{H} = c_1 - \frac{N_s \ddot{H}}{A \dot{\theta}}$$

$$\mathcal{Y}'' + \left[ \left( \frac{4}{r} + \Phi' - \Lambda' \right) + \frac{\mu'}{\mu} \right] \mathcal{Y}' + \left[ \frac{\epsilon + p}{\mu} \omega^2 e^{-2\Phi} - \frac{(\ell + 2)(\ell - 1)}{r^2} \right] e^{2\Lambda} \mathcal{Y} = 0.$$

# identification of SGR 1806-20

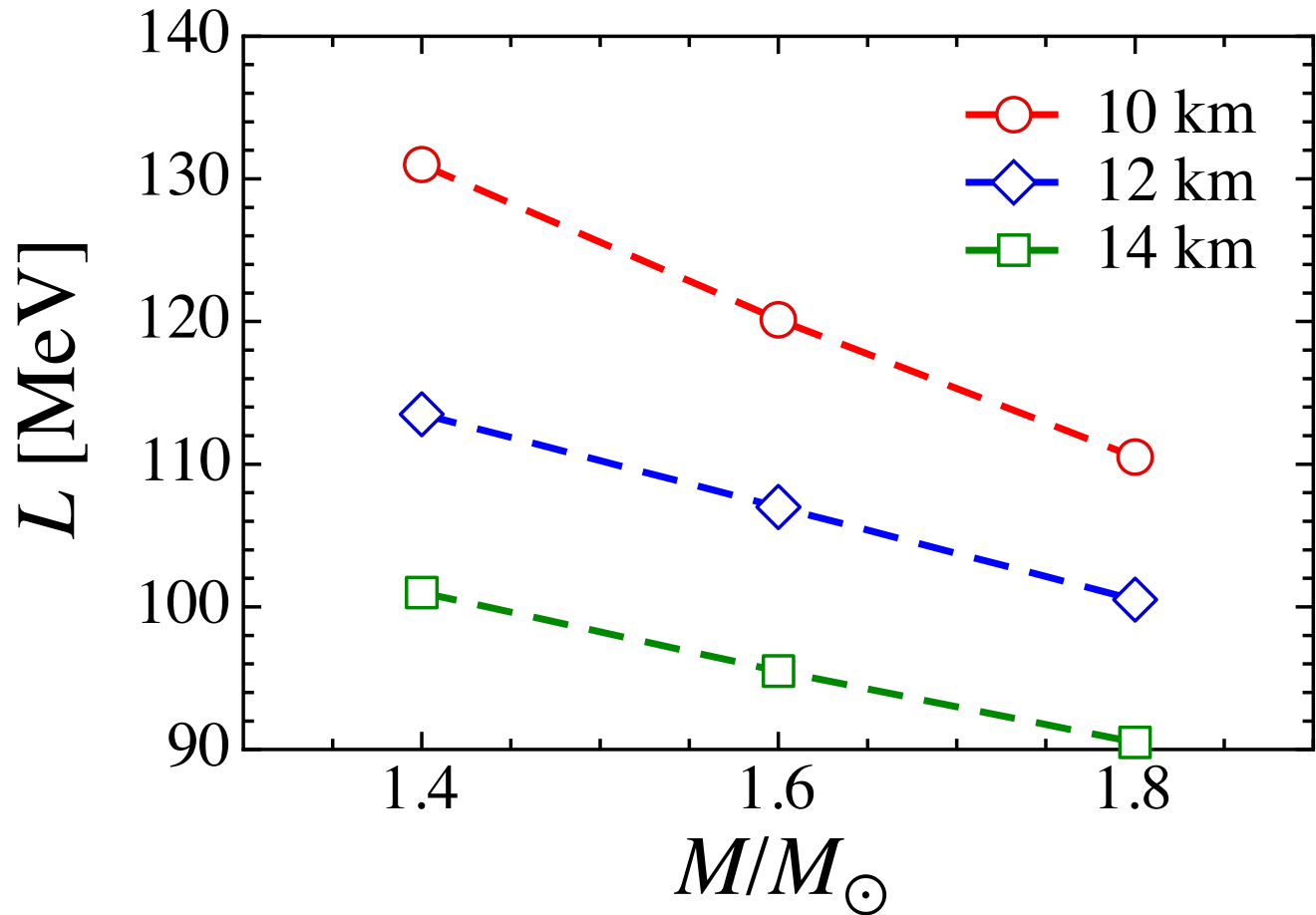


# constraint on $L$ via SGR 1806-20



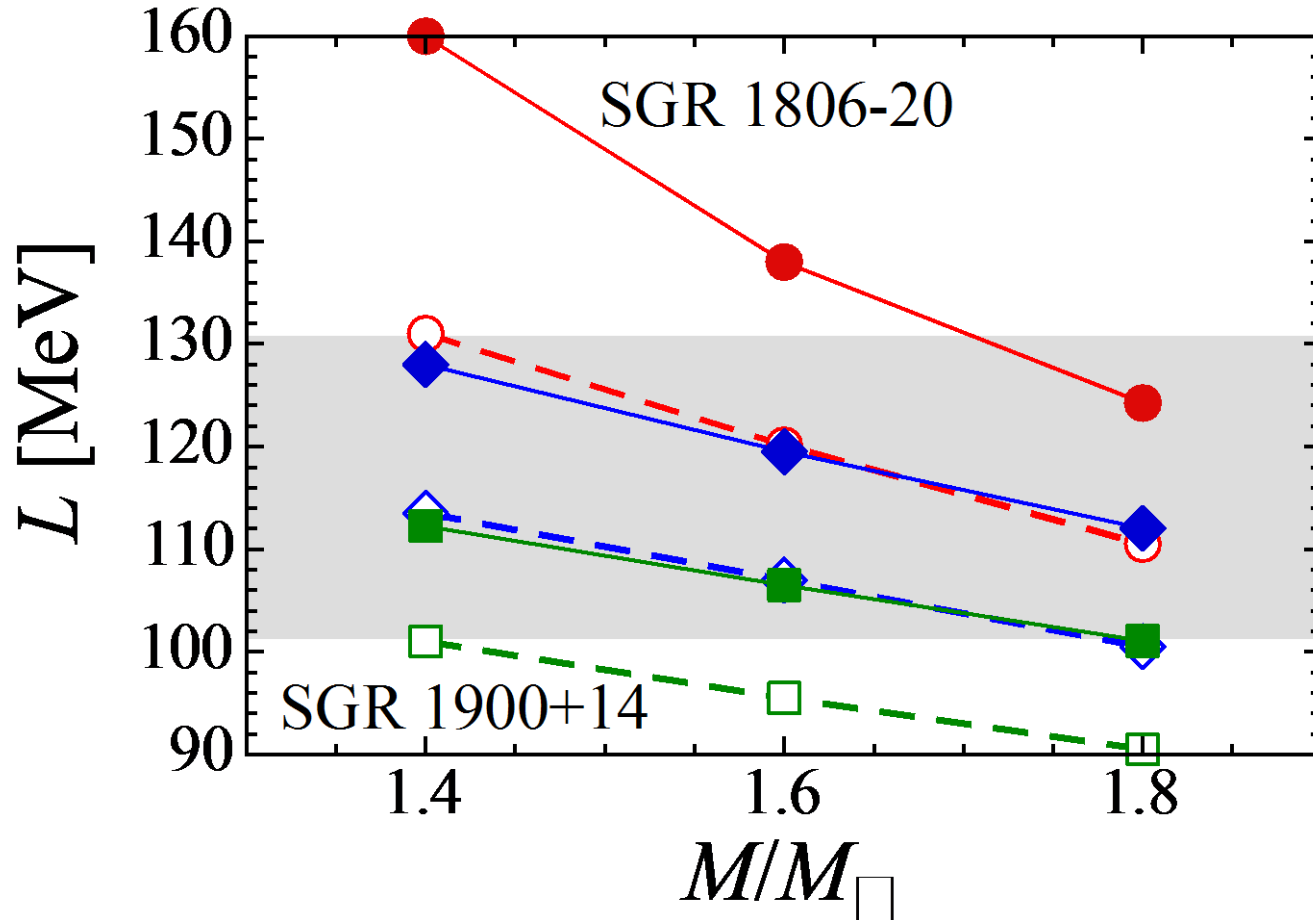
➔  $101.1 \text{ MeV} \leq L \leq 160.0 \text{ MeV}$

# constraint on $L$ via SGR 1900+14



➔  $90.5 \text{ MeV} \leq L \leq 131.0 \text{ MeV}$

# allowed region for $L$



→  $101.1 \text{ MeV} \leq L \leq 131.0 \text{ MeV}$

# missing effects ??

blue : decrease

red : increase

- modification of shear modulus
  - size of nuclei
  - electron screening (Horowitz & Hughto 08; Kobayakov & Pethick 13)
  - existence of pasta phase (Sotani 11; Gearheart+11; Newton+13)
- pairing effect and shell effect (Deibel+13)
- superfluidity (Chamel 12, 13; Sotani+12; Deibel+13)
- magnetic field (Sotani+; Colaiuda & Kokkotas; Gabler+; Passamonti+; Lander+; Deibel+13)
- emission mechanism ??

# summary

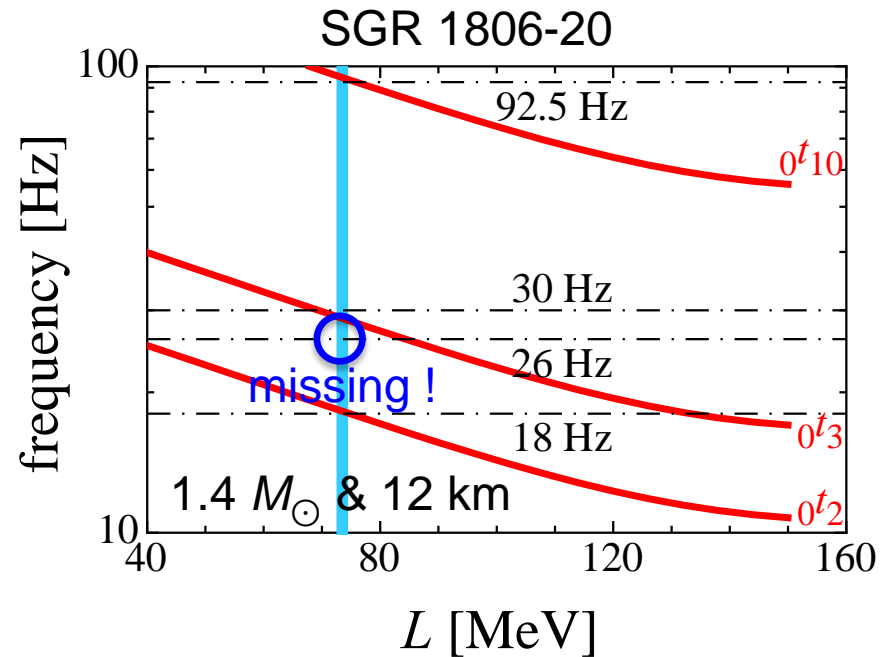
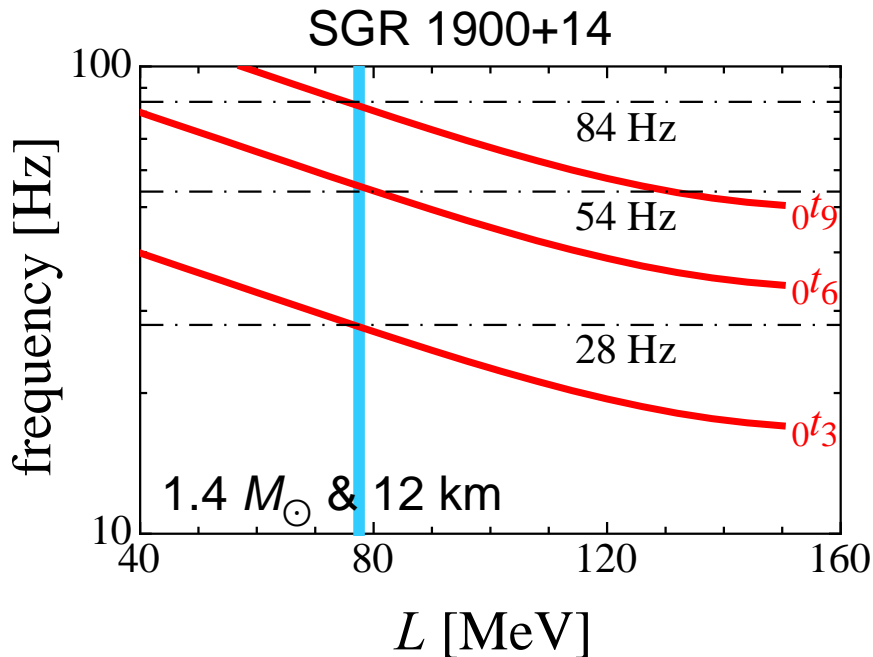
- **asteroseismology** could be powerful approach to see the interior properties of neutron stars.
  - QPOs in SGRs may be good examples to adopt the asteroseismology
- comparing the torsional oscillations to the observational evidences, we can get the constraint on  $L$  as  $L \gtrsim 50 \text{ MeV}$ .
- superfluid effect enhances the frequencies of torsional oscillations.
  - $100 \lesssim L \lesssim 130 \text{ MeV}$ , if all QPOs come from torsional oscillations
  - $58 \lesssim L \lesssim 85 \text{ MeV}$ , if QPOs except for 26 Hz QPO come from torsional oscillations
- we should take into account additional missing effects.





# alternative possibility

instead of previous correspondence, i.e.,  $l = 4, 8, 13$  for SGR 1900+14, and  $l = 3, 4, 5, 15$  for SGR 1806-20, we may consider alternative possibility as



26 Hz QPO observed in SGR 1806-20 remains a complete puzzle !!

# relative error

- previous identification

QPOs (Hz)	$l$	${}_0t_l$ (Hz)	error (%)
18	3	18.50	-2.79
26	4	24.82	4.53
30	5	30.96	-3.19
92.5	15	90.18	2.51

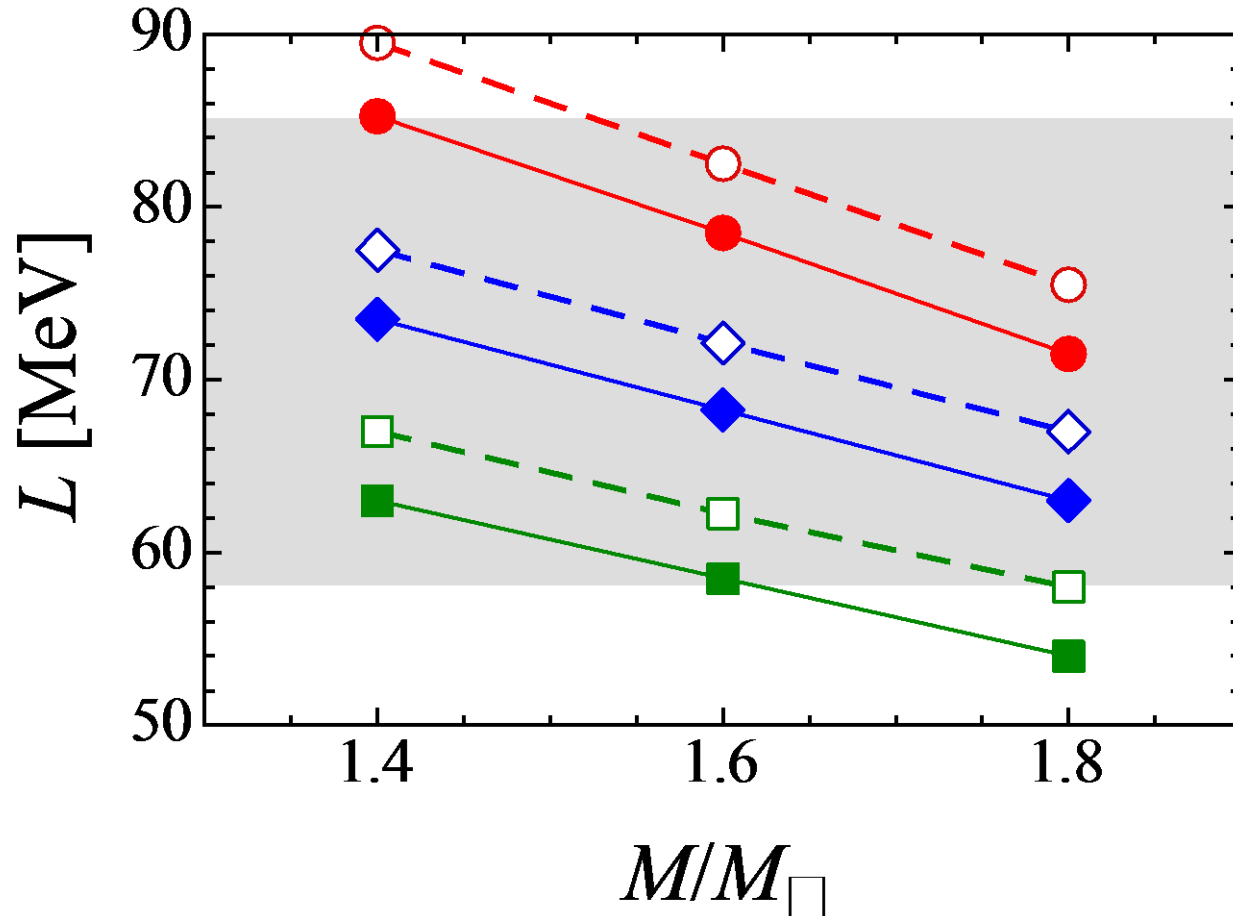
QPOs (Hz)	$l$	${}_0t_l$ (Hz)	error (%)
28	4	27.26	2.63
54	8	53.76	4.50
84	13	86.18	-2.60

- alternative identification

QPOs (Hz)	$l$	${}_0t_l$ (Hz)	error (%)
18	2	18.23	-1.27
26	---	---	---
30	3	28.82	3.93
92.5	10	94.70	-2.38

QPOs (Hz)	$l$	${}_0t_l$ (Hz)	error (%)
28	3	27.74	0.93
54	6	55.48	-2.74
84	9	82.29	2.04

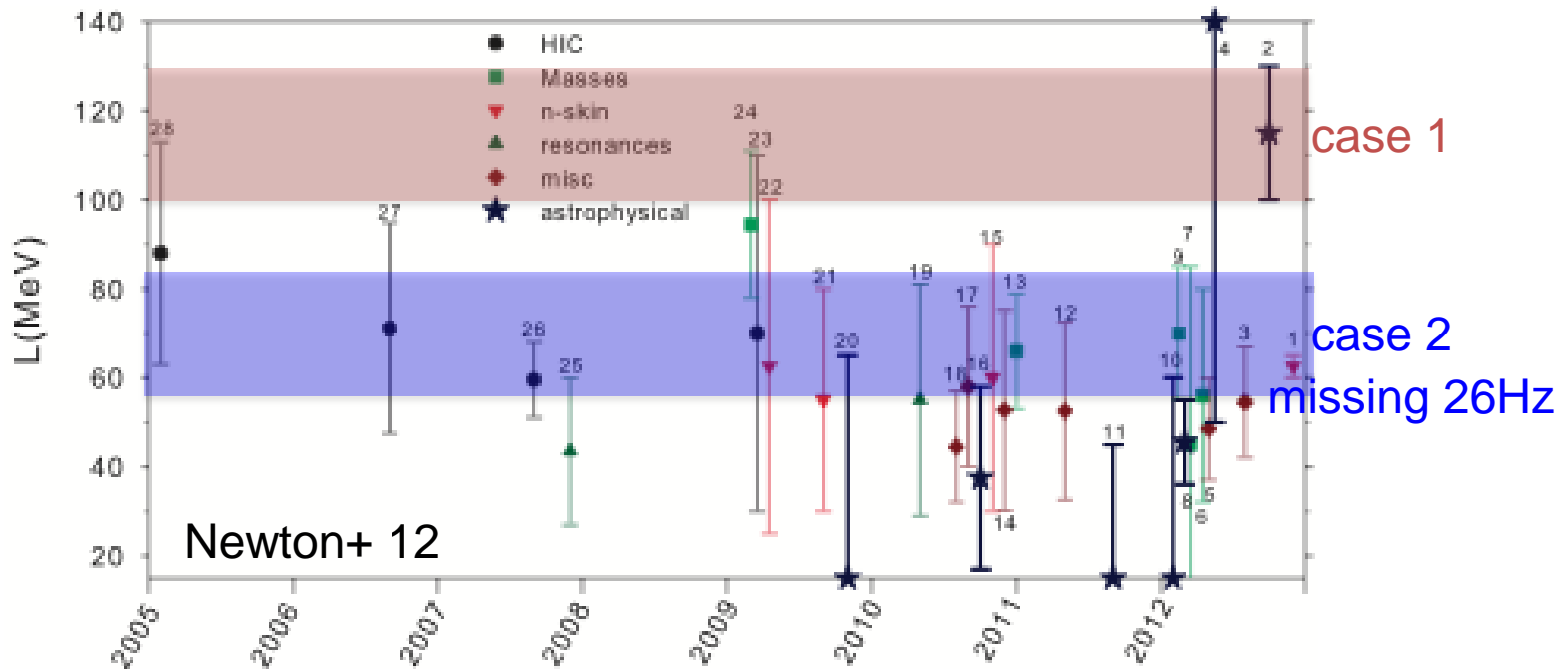
# alternative allowed region for $L$



→  $58.0 \text{ MeV} \leq L \leq 85.3 \text{ MeV}$   
( $32.4 \text{ MeV} \leq S_0 \leq 34.4 \text{ MeV}$ )

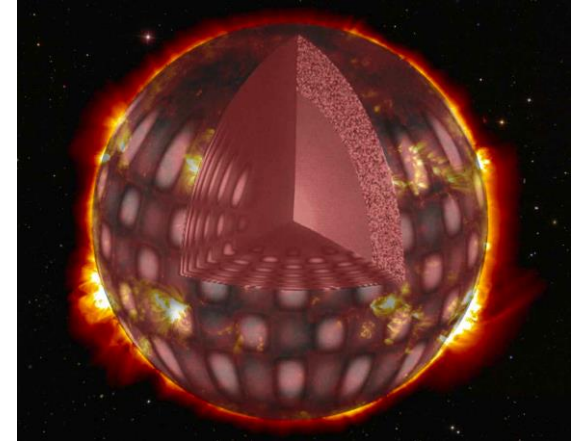
# other constraints on $L$

- other constraints suggests  $L \sim 60 \pm 20$  MeV ?
  - this means case 2 may be favored ??
  - if so, one has to prepare another oscillation mechanism...



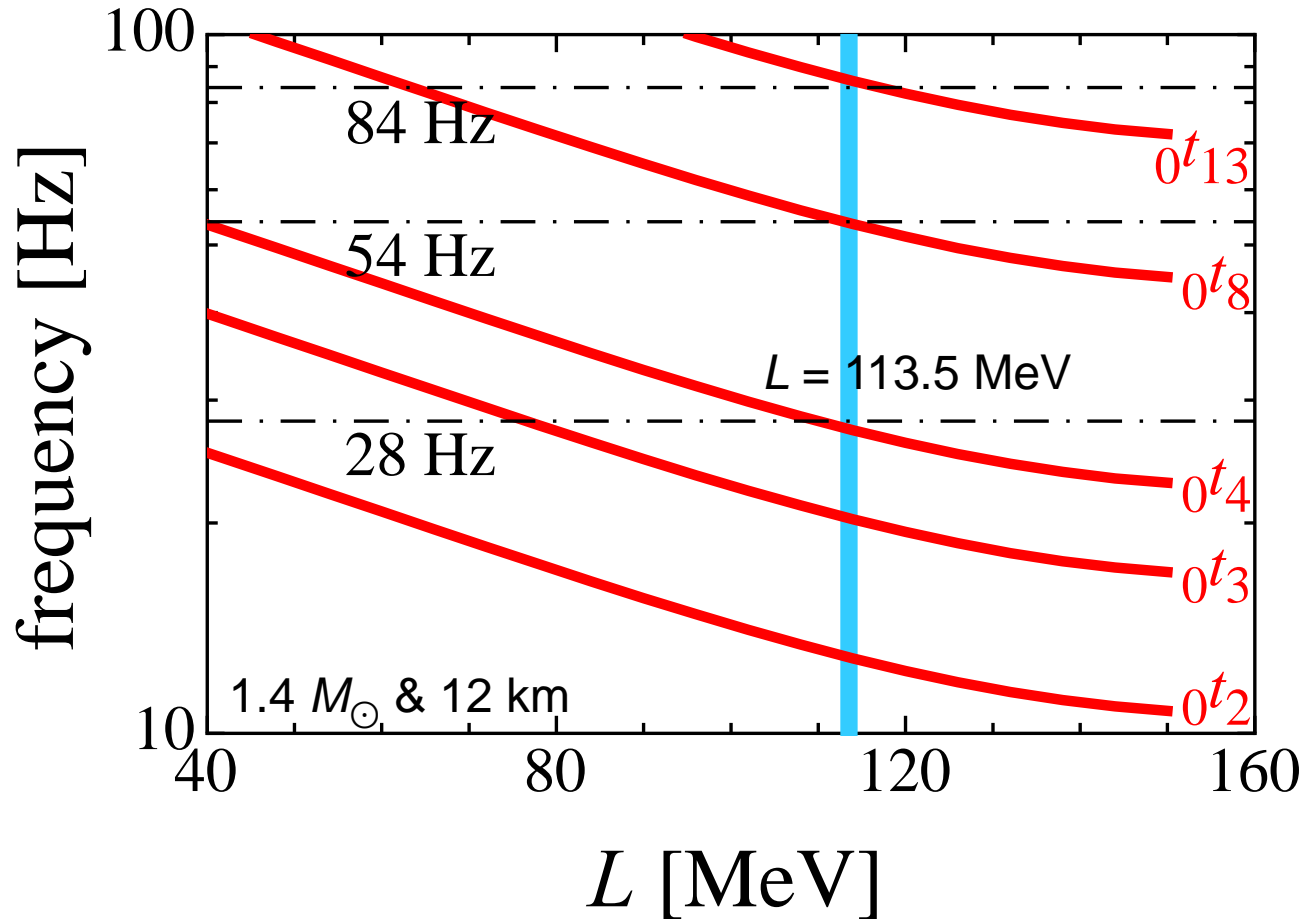
# oscillations in NSs

- polar oscillations
  - fluid modes
    - ✓ fundamental mode (f -mode) ... ~ kHz
    - ✓ pressure mode (p-mode) ... > a few kHz
    - ✓ gravity mode (g-mode) ... < a few 100 Hz
    - ✓ rotational mode (r-mode) ... ~ rotation frequency
  - relativistic modes
    - ✓ spacetime mode (w-mode) ... > a few tens kHz
- axial oscillations
  - fluid modes; **torsional mode (t-mode)** ... > ten Hz
  - relativistic modes; w-mode ... > a few tens kHz

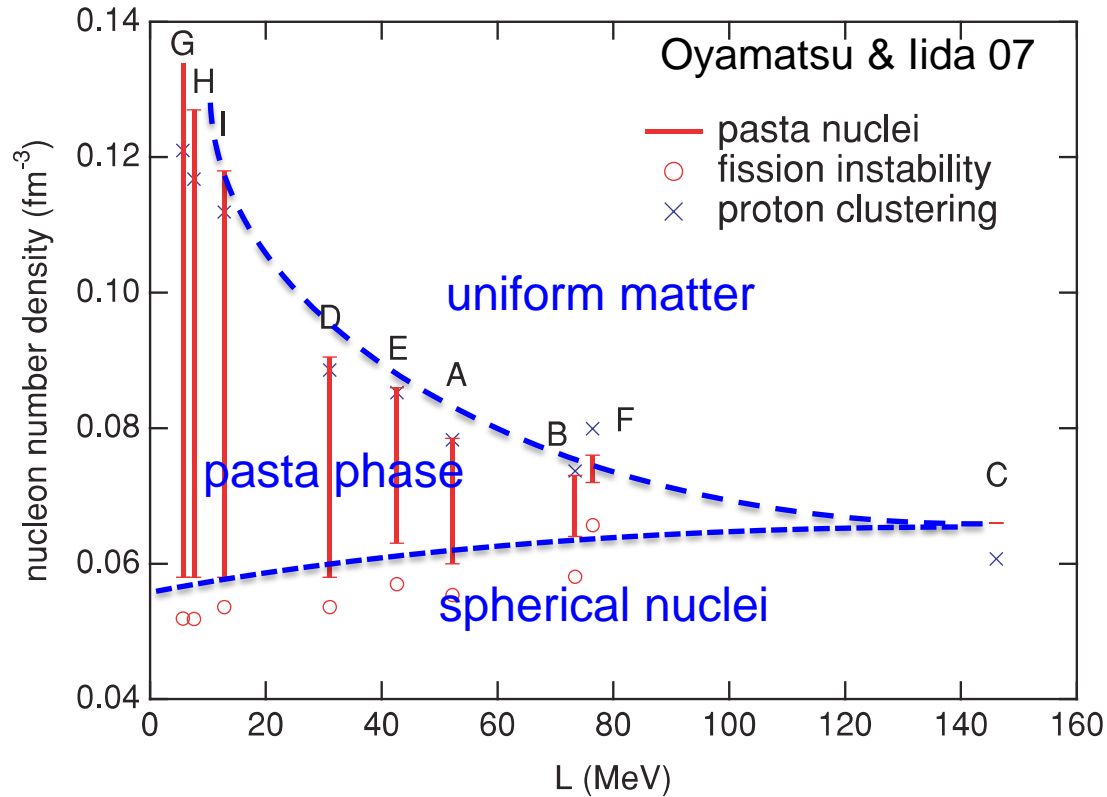


©NASA

# identification of SGR 1900+14



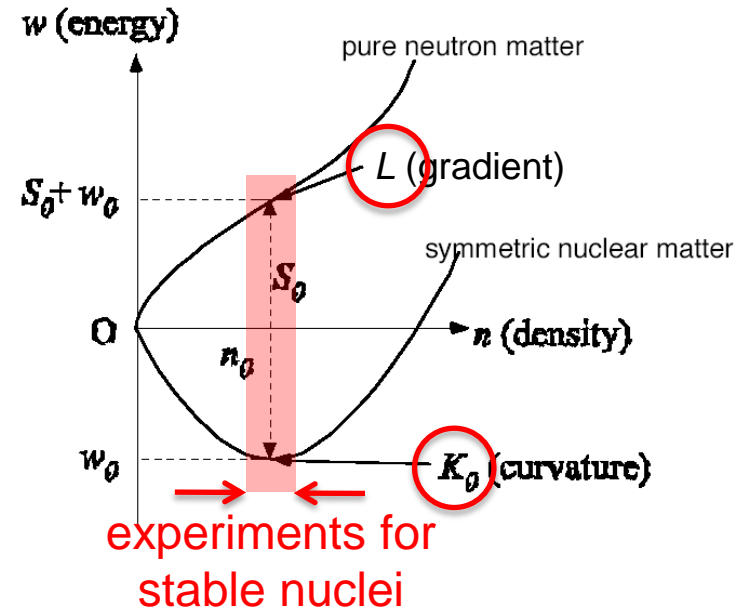
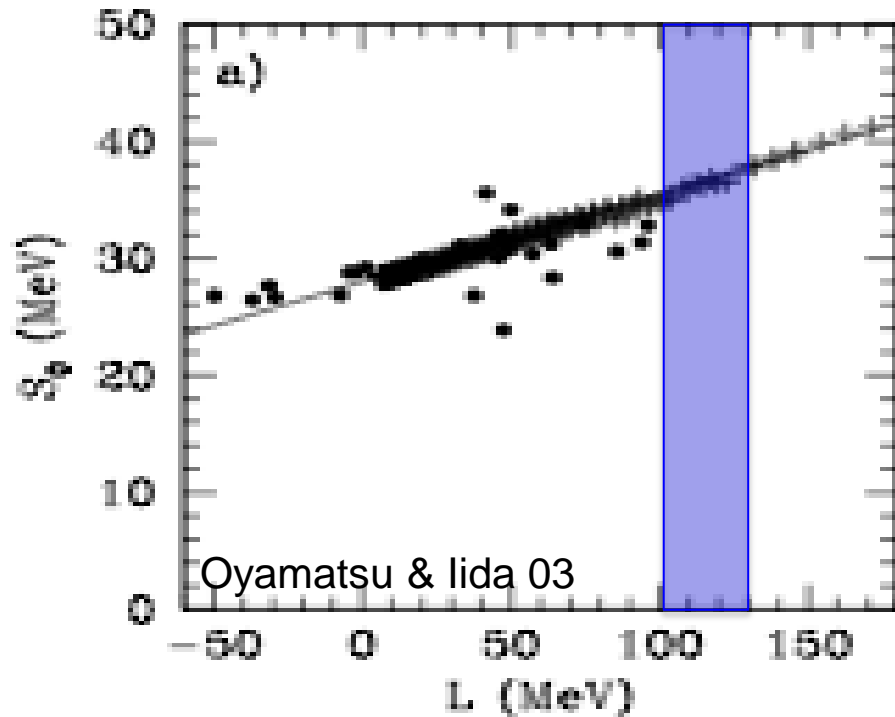
# pasta phase



- region of pasta phase depends strongly on  $L$
- for  $L \gtrsim 100\text{MeV}$ , pasta structure almost disappears

# constraint on $S_0$

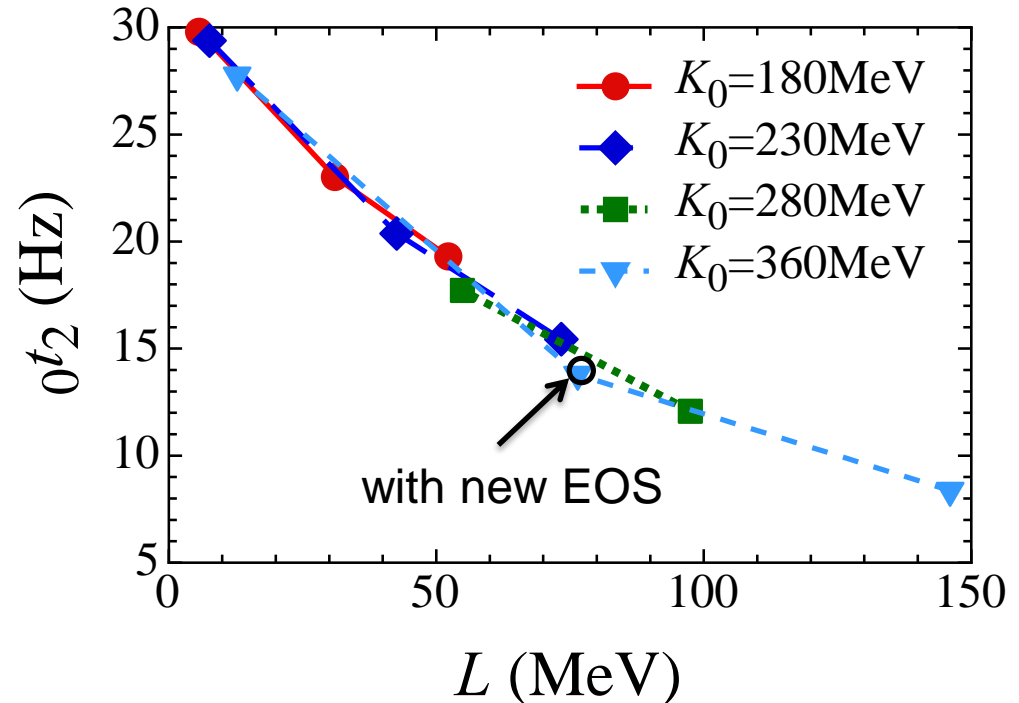
- by using the empirical relation :  $S_0 = 28 + 0.075L$   
 (Oyamatsu & Iida 03)
- $35.6 \text{ MeV} \leq S_0 \leq 37.8 \text{ MeV}$





# comparing with other EOS

- new EOS (Miyatsu+ 13)
  - core : RMF calculation
  - crust : TF theory
- EOS parameters
  - $L = 77.1$  MeV
  - $K_0 = 274$  MeV
- even with new EOS, the dependence of torsional oscillations on  $L$  is same as the previous results.



# effect of electron screening

- contribution due to Coulomb interaction
  - Ogata & Ichimaru 90; Strohmayer+ 91

$$\mu = 0.1194 \times \frac{n_i (Ze)^2}{a}$$

- including effect of electron screening
  - Horowitz & Hughto 08 : 10% reduction
  - Kobayakov & Pethick 13

$$\mu = 0.1194 \left[ 1 - 0.010 Z^{2/3} \right] \frac{n_i (Ze)^2}{a}$$

effect of electron screening

- ~11.7% reduction for  $Z = 40$
- phonon contribution is much smaller (Baiko 12)

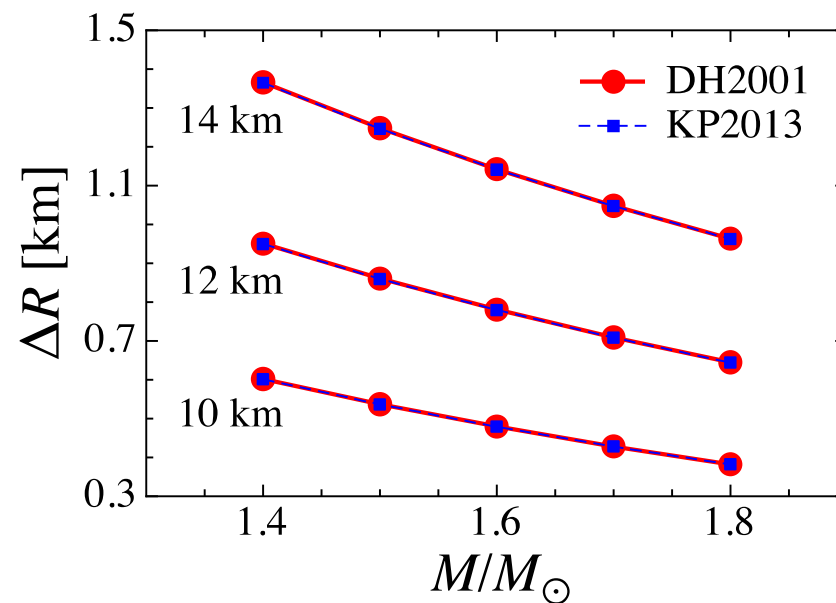
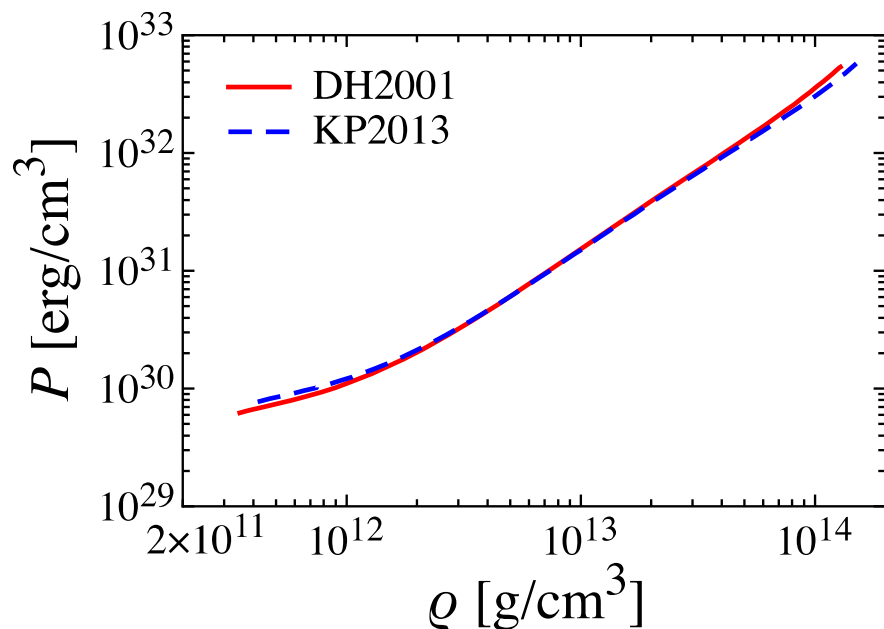
# adopted EOS

- outer crust
  - Haensel & Pichon 94
- inner crust
  - compressible liquid drop model (CLDM)
  - Kobyakov & Pethick 13 (KP2013) based on Lattimer & Swesty 91
  - Douchin & Haensel 01 (DH2001)

	KP2013	DH2001
model	CLDM	CLDM
neutron skin	×	○
$n_{bc}$ [1/fm <sup>3</sup> ]	$8.913 \times 10^{-2}$	$7.596 \times 10^{-2}$
$\rho_c$ [g/cm <sup>3</sup> ]	$1.504 \times 10^{14}$	$1.285 \times 10^{14}$

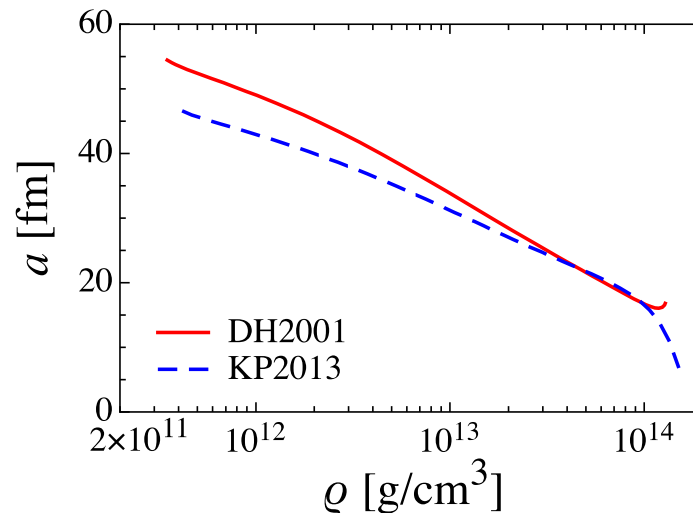
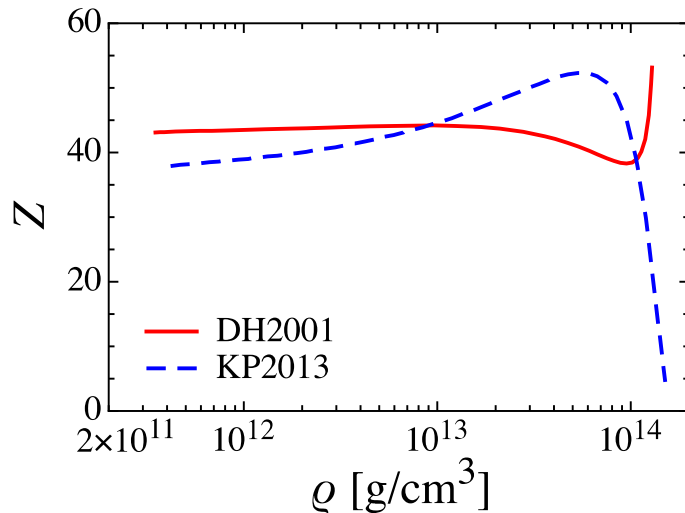
# crust models

- quite difficult to distinguish the difference in crust thickness with DH2001 and with KP2013



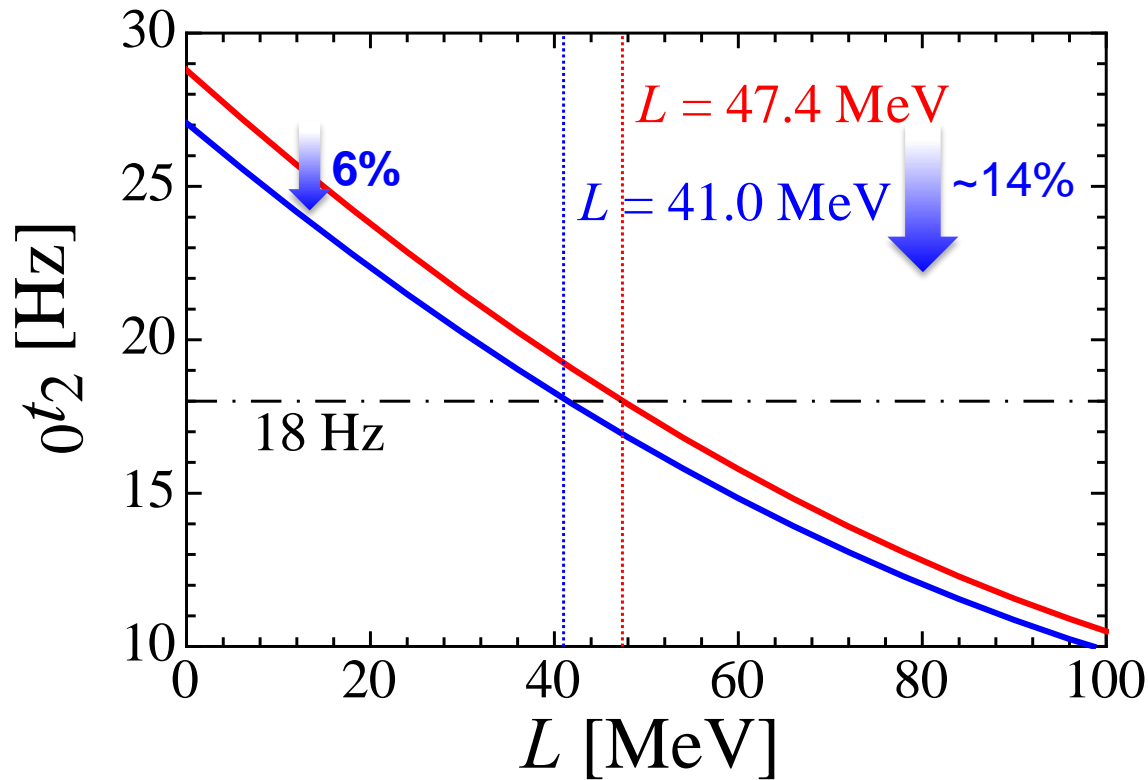
# other properties

- charge depends strongly on the EOS at the crust basis
- radius of WS cell also depends on the EOS



# constraint on $L$

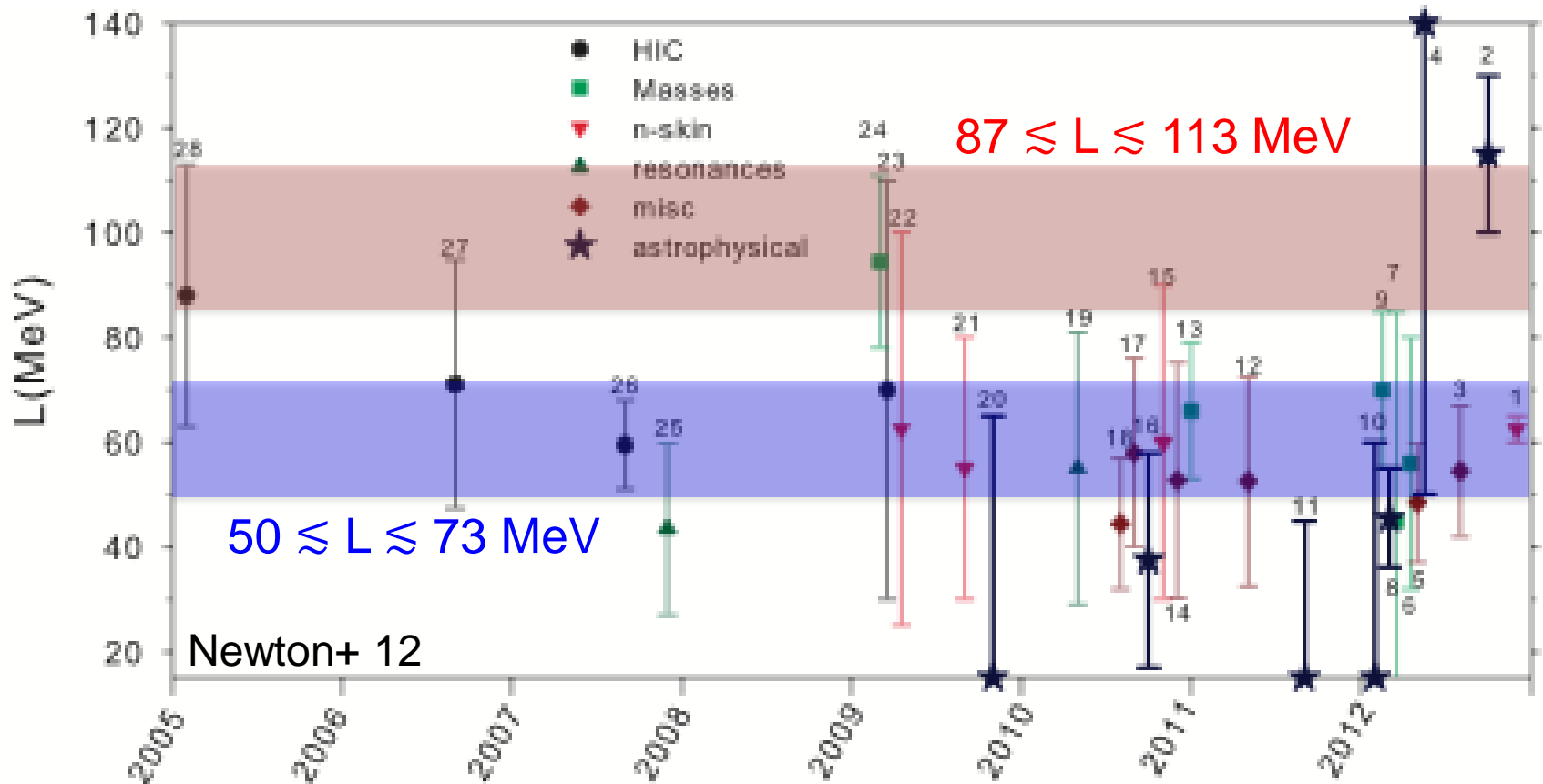
- due to the electron screening effect, constraint of  $L$  shifts ~14% smaller value



$L \gtrsim 47.4$  MeV  $\longrightarrow$   $L \gtrsim 41.0$  MeV

# modified constraints on $L$

- adopting the reduction of frequencies due to the electron screening effect, constraints on  $L$  become as follows;



# fundamental oscillations

(HS 13)

- one may be identify the EOS using the observations of crustal oscillations
- independent of the stellar mass and the crust EOS, the effect of electron screening can reduce 6% of the frequencies

