Exact solution of Einstein's angular momentum constraint in the angular-momentum sector of matter on the past light cone:

Exact dragging of inertial axes by cosmic energy currents (Mach's principle)

Christoph Schmid, ETH Zurich

27th Texas Symposium on Relativistic Astrophysics December 8-13, 2013, Dallas, Texas

Inertial Axes \Leftrightarrow Spin Axes of Gyroscopes

 Experimentally spin axes of gyroscopes directly give time-evolution of local inertial axes: inertial guidance systems.

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

Inertial Axes \Leftrightarrow Spin Axes of Gyroscopes

 Experimentally spin axes of gyroscopes directly give time-evolution of local inertial axes: inertial guidance systems.

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

 Argument in opposite direction: In the local inertial frame of a gyroscope: no gyroscope-precession.

Super-precise observational fact

- Spin-axes of gyroscopes do not rotate relative to quasars, except extremely small dragging effect by Earth rotation:
 Lense - Thirring effect, detected by Gravity Probe B.
- Gyros at a few Earth radii: Dragging by Earth negligible.

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

Mach's Principle: The Question

WHAT PHYSICAL CAUSE:

determines the **time-evolution of gyroscope axes**, i.e. the **time-evolution of inertial axes** ?

▲□▶▲□▶▲□▶▲□▶ □ のQ@

Mach's Principle: The Question

WHAT PHYSICAL CAUSE:

determines the **time-evolution of gyroscope axes**, i.e. the **time-evolution of inertial axes** ?

John A. Wheeler: "Who gives the marching orders" to gyroscope axes (= inertial axes)?

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

The Postulate formulated by Mach

Inertial axes exactly follow an average of the motion of cosmological masses: exact frame dragging.

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

The Postulate formulated by Mach

- Inertial axes exactly follow an average of the motion of cosmological masses: exact frame dragging.
- Newton's gravitational force: no torque on a gyroscope. Mach wrote: unknown, what new force could do the job. General Relativity: gravito-magnetism

 Lense-Thirring.

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

The Postulate formulated by Mach

- Inertial axes exactly follow an average of the motion of cosmological masses: exact frame dragging.
- Newton's gravitational force: no torque on a gyroscope.
 Mach wrote: unknown, what new force could do the job.
 General Relativity: gravito-magnetism
 Lense-Thirring.
- Mach did not know, what average of cosmological masses and their motions should be taken.

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

Mach's Principle: Our Results

• Our old results, published in Phys. Rev. 2006 and 2009:

exact dragging of inertial axes (Mach's Principle) for all linear perturbations on all FRW backgrounds.

(ロ) (同) (三) (三) (三) (三) (○) (○)

Mach's Principle: Our Results

• Our old results, published in Phys. Rev. 2006 and 2009:

exact dragging of inertial axes (Mach's Principle) for all linear perturbations on all FRW backgrounds.

Our new results, unpublished:

Exact angular momentum constraint equation for matter angular-momentum on past light-cone: EXACT Einstein equation: LINEAR ord. diff. eq.

Mach's postulate \Rightarrow LINEAR relation input, output.

Significance: Is rotation absolute or relative?

- LOCAL analysis in laboratory, solar system, galaxy:
- Analysis in solar system: Newton correct with ABSOLUTE-UNAMBIGOUS non-rotating frame.
- General Relativity of solar system (e.g. perihelion shift): Local non-rotating frames:

ABSOLUTE-UNAMBIGUOUS,

Asymptotic **nonrotating frame**: **observational input**, theoretically:

ABSOLUTE-NO-CAUSE-input-by-hand

via spatial boundary condition,

GR in solar system is not a "Theory of Relativity".

(ロ) (同) (三) (三) (三) (三) (○) (○)

Significance: Is rotation absolute or relative?

- LOCAL analysis in laboratory, solar system, galaxy:
- Analysis in solar system: Newton correct with ABSOLUTE-UNAMBIGOUS non-rotating frame.
- General Relativity of solar system (e.g. perihelion shift): Local non-rotating frames:

ABSOLUTE-UNAMBIGUOUS,

Asymptotic **nonrotating frame**: **observational input**, theoretically:

ABSOLUTE-NO-CAUSE-input-by-hand

via spatial boundary condition,

GR in solar system is not a "Theory of Relativity".

(ロ) (同) (三) (三) (三) (三) (○) (○)

► GLOBAL COSMOLOGICAL analysis:

 Local non-rotating system is nonrotating RELATIVE to average of cosmic matter: Cosmological GR is a "Theory of Relativity".

NON-PERTURBATIVE cosmological geometry

- Spirit of **Riemann normal coord.** on arbitrary geometry.
- Observation event P₀: World-lines of photons observed give radial coordinate lines on past light-cone (LC).

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

► Observer at rest relative to asymptotic quasars ⇔ no dipole term in measured Hubble velocities.

NON-PERTURBATIVE cosmological geometry

- Spirit of Riemann normal coord. on arbitrary geometry.
- Observation event P₀: World-lines of photons observed give radial coordinate lines on past light-cone (LC).
- Observer at rest relative to asymptotic quasars

 no dipole term in measured Hubble velocities.
- Minkowski corridor along photon world-line arriving at P₀: Local Minkowski Coordinate System at P₀ extended along photon world-line gives spatial separation in frame of observer (astronomer's luminosity distance),

 $P \Rightarrow \text{coordinate } r_P \text{ with } g_{rr} = 1.$

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

NON-PERTURBATIVE cosmological geometry

- Spirit of **Riemann normal coord.** on arbitrary geometry.
- Observation event P₀: World-lines of photons observed give radial coordinate lines on past light-cone (LC).
- ► Observer at rest relative to asymptotic quasars ⇔ no dipole term in measured Hubble velocities.
- Minkowski corridor along photon world-line arriving at P₀: Local Minkowski Coordinate System at P₀ extended along photon world-line gives spatial separation in frame of observer (astronomer's luminosity distance),

 $P \Rightarrow$ coordinate r_P with $g_{rr} = 1$. \triangleright Choose polar quasar and zero-meridian quasar

 $P \Rightarrow \text{ coordinates } (\theta_P, \phi_P).$

► Equilateral triangles, infinitely narrow: two neighboring photon world-lines with equal spatial distances $r_{OP} = r_{OP'}$ generate **2-spheres orthogonal** to photon world-lines:

$$g_{r\theta}=0, \qquad g_{r\phi}=0.$$

Gyroscope precession

⇒ toroidal vorticity sector for cosmic matter flow

- Non-perturbative treatment. Exact fields.
- On one past light-cone: decomposition in 3-scalar, 3-vector, 3-tensor sectors rigorously valid.
- On one past light-cone: We neglect tensor sector.
- Scalar sector ⇒ no gyroscope-precession ⇒ neglect it.

(ロ) (同) (三) (三) (三) (三) (○) (○)

Gyroscope precession

⇒ toroidal vorticity sector for cosmic matter flow

- Non-perturbative treatment. Exact fields.
- On one past light-cone: decomposition in 3-scalar, 3-vector, 3-tensor sectors rigorously valid.
- On one past light-cone: We neglect tensor sector.
- Scalar sector ⇒ no gyroscope-precession ⇒ neglect it.
- Gyroscope precession at P_O caused only by:
 vector sector (divergence-free) ≡ vorticity sector
 with J^P = 1⁺ ⇒ toroidal vorticity, relative to P_O.

(日) (日) (日) (日) (日) (日) (日)

▶ focus on m = 0. With $m = \pm 1$ following trivially.

Gyroscope precession

⇒ toroidal vorticity sector for cosmic matter flow

- Non-perturbative treatment. Exact fields.
- On one past light-cone: decomposition in 3-scalar, 3-vector, 3-tensor sectors rigorously valid.
- On one past light-cone: We neglect tensor sector.
- Scalar sector ⇒ no gyroscope-precession ⇒ neglect it.
- Gyroscope precession at P_O caused only by:
 vector sector (divergence-free) ≡ vorticity sector
 with J^P = 1⁺ ⇒ toroidal vorticity, relative to P_O.
- ▶ focus on m = 0. With $m = \pm 1$ following trivially.
- Vector spher. harmonics, toroidal sector:

$$(V_{\ell=1, m=0}^{(P=+)})^{\phi} = \text{constant}, \qquad (V_{\ell=1, m=0}^{(P=+)})^{\theta} = 0,$$

give angular velocity of rigid rotation around z-axis.

Light cones: Intrinsic geometry UNPERTURBED by strong toroidal vorticity

• Measure circumference C_P of 2-sphere through P. For simplicity in this talk, assume that $(C_P/r_P) = 2\pi \Rightarrow$ measured light-cone indistinguishable from light cones in Minkowski and in spatially flat FRW.

(ロ) (同) (三) (三) (三) (三) (○) (○)

Light cones: Intrinsic geometry UNPERTURBED by strong toroidal vorticity

- Measure circumference C_P of 2-sphere through P. For simplicity in this talk, assume that $(C_P/r_P) = 2\pi \Rightarrow$ measured light-cone indistinguishable from light cones in Minkowski and in spatially flat FRW.
- Tensor spherical harmonics with toroidal vorticity,

$$(T_{\ell=1, m=0}^{(P=+)})_{\alpha\beta} = 0.$$

Metric of 2-spheres must be unperturbed metric,

$$g_{\theta\phi} = 0, \quad g_{\theta\theta} = r^2, \quad g_{\phi\phi} = r^2 \sin^2 \theta.$$

INTRINSIC 3-GEOMETRY of every light-cone (LC) with ARBITRARY VORTICITY fields:

 \Rightarrow EQUAL to UNPERTURBED intrinsic geometry of LC in Minkowski and LC in spatially flat FRW.

・ロト・日本・日本・日本・日本

From one light cone to a neighbouring one: SHIFT

- Lapse function (elapsed measured time between slices at fixed χ) is unperturbed, because it is a 3-scalar.
- Only quantity referring to toroidal vorticity: shift-3-vector

$$\ell = 1, m = 0: \quad \beta^{\phi} \equiv \beta(\mathbf{v}, \chi), \quad \beta^{\theta} = 0, \quad \beta^{\chi} = 0.$$

(ロ) (同) (三) (三) (三) (三) (○) (○)

From one light cone to a neighbouring one: SHIFT

- Lapse function (elapsed measured time between slices at fixed χ) is unperturbed, because it is a 3-scalar.
- Only quantity referring to toroidal vorticity: shift-3-vector

$$\ell = 1, m = 0: \quad \beta^{\phi} \equiv \beta(\mathbf{v}, \chi), \quad \beta^{\theta} = 0, \quad \beta^{\chi} = 0.$$

• Hubble H(r) measured

comoving coordinate difference $\equiv d\chi \equiv a^{-1} dr$,

at fixed comoving distance: measured time diff. $\equiv dt$,

conformal time difference $\equiv d\eta \equiv a^{-1}dt$.

• Light-cone coordinates (v, χ, θ, ϕ) and metric:

 $\begin{array}{lll} \mathbf{v} &\equiv & \eta + \chi & \Rightarrow & \mathbf{v} \text{ labels past light-cone,} \\ \mathbf{ds}^2 &= & \mathbf{a}^2 [-\mathbf{dv}^2 + 2\mathbf{dv} \, \mathbf{d}\chi + \chi^2 (\mathbf{d}\theta^2 + \sin^2\theta \, \mathbf{d}\phi^2) \\ &+ 2 \,\beta_\phi \, \chi^2 \sin^2\theta \, \mathbf{d}\phi \, \mathbf{d}v]. \end{array}$

Exact angular momentum constraint for toroidal vorticity on past light-cone of observation

• Exact Einstein tensor for vorticity with $(\ell = 1, m = 0)$,

$$egin{aligned} &\mathbf{2}a^{4}G_{(1)}^{
u\phi} = \ &= \partial_{\chi}^{2}eta + (\partial_{\chi}eta)\left(4\chi^{-1} - 2\mathcal{H}
ight) + eta\left(4\mathcal{H}' - 4\mathcal{H}^{2}
ight), \end{aligned}$$

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

where $\mathcal{H} \equiv a^{-1}(da/d\eta)$ and $\mathcal{H}' \equiv (d\mathcal{H}/d\eta)$.

Exact angular momentum constraint for toroidal vorticity on past light-cone of observation

• Exact Einstein tensor for vorticity with $(\ell = 1, m = 0)$,

$$\begin{split} & 2a^4 G^{\nu\phi}_{(1)} = \\ & = \partial_{\chi}^2 \beta + (\partial_{\chi}\beta) \left(4\chi^{-1} - 2\mathcal{H} \right) + \beta \left(4\mathcal{H}' - 4\mathcal{H}^2 \right), \end{split}$$

where $\mathcal{H} \equiv a^{-1}(da/d\eta)$ and $\mathcal{H}' \equiv (d\mathcal{H}/d\eta)$.

▶ Matter evolution: If p/ρ approximated time-independent, $a_0 \equiv 1, \ \eta_0 \equiv 1$: $a = \eta^P \implies \mathcal{H} = P \eta^{-1}$. Equation for Green function away from δ -function shell,

$$\begin{split} \partial_{\eta}^{2}\beta &- (\partial_{\eta}\beta) \left[4(1-\eta)^{-1} - 2P\eta^{-1} \right] \\ &- 4\beta \left(P + P^{2} \right) \eta^{-2} = 0, \end{split}$$

from $\eta = 0$ (big bang) to $\eta = 1$ (observation): Singular points of regular type, solutions with algebraic singularities. **Numerical solution** of linear diff. eq. of 2nd order needed.

General Relativity \Rightarrow Gravitomagnetism

- Work on Uniform-Radial-Hubble-Expansion Slices
- Operational definitions via measurements by Fiducial Observers (FIDOs) at rest on slice with Local Ortho-Normal Bases (LONBs) denoted by hats.
- Op. Def. of gravitoelectric field $\vec{E}_{\rm g} \equiv \vec{g}$:

 $\frac{d}{dt}p_{\hat{j}} \equiv m E_{\hat{j}}^{g}$ free-falling **quasistatic** test particle,

• Op. Def. of gravitomagnetic field \vec{B}_{g} :

$$\Omega_{\hat{i}}^{\text{gyro}} \equiv -\frac{1}{2}B_{\hat{i}}^{\text{g}}$$
 precession of gyro comoving with FIDO.

► Gravitomagnetic vector potential \vec{A}_{g} : \vec{B}_{g} = curl \vec{A}_{g} it follows : \vec{A}_{g} = $\vec{\beta}$ ≡ shift vector.

Einstein's $G_{\hat{k}}^{\hat{o}}$ Eq.: Angular Momentum Constraint

 NEW: Angular momentum constraints for Uniform-Radial-Expansion Slices are form-identical for slices which are spatially flat, hyperbolic, or spherical,

$$(-\Delta + \mu^2) \, \vec{A}_{\mathrm{g}} = -16\pi G_{\mathrm{N}} \, \vec{J}_{\varepsilon},$$

$$(\mu/2)^2 \equiv -(dH/dt) \equiv (H$$
-dot radius)⁻²,
 $\vec{J_{\varepsilon}} =$ energy current density = momentum density

NEW: Same as Ampère's law for stationary magnetism except µ² Å_g ⇒ Yukawa suppression (*H*-dot radius).
 Elliptic equation (no partial time-derivatives of perturbations) for time-dependent gravitomagnetism.

Einstein's $G_{\hat{k}}^{\hat{o}}$ Eq.: Angular Momentum Constraint

 NEW: Angular momentum constraints for Uniform-Radial-Expansion Slices are form-identical for slices which are spatially flat, hyperbolic, or spherical,

$$(-\Delta + \mu^2) \, \vec{A}_{\mathrm{g}} = -16\pi G_{\mathrm{N}} \, \vec{J}_{\varepsilon},$$

$$(\mu/2)^2 \equiv -(dH/dt) \equiv (H$$
-dot radius)⁻²,
 $\vec{J}_{\varepsilon} =$ energy current density = momentum density

- ► NEW: Same as Ampère's law for stationary magnetism except µ² A_g ⇒ Yukawa suppression (*H*-dot radius). Elliptic equation (no partial time-derivatives of perturbations) for time-dependent gravitomagnetism.
- ► NEW: $\vec{J}_{\hat{k}}^{\varepsilon} = T_{\hat{k}}^{\hat{0}}$ is measurable input without prior knowledge of g_{0i} , which is output. T_k^0 not measurable.
- NEW in cosmology: Laplacian on Vector Fields in Riemannian 3-Spaces is de Rham - Hodge Laplacian.

シック・ 川 ・ 山 ・ 小田 ・ 小田 ・ 小田 ・

Solution of the Angular Momentum Constraint

 Identical expressions for solutions of cosmological gravitomagnetism on Uniform-Radial-Expansion slices, which are spatially open,

$$\vec{B}_{g}(P) = -2 \vec{\Omega}_{gyro}(P) =$$

$$= -4 G_{N} \int d(vol_{Q}) [\vec{n}_{PQ} \times \vec{J}_{\varepsilon}(Q)] Y_{\mu}(r_{PQ})$$

$$Y_{\mu}(r) = \frac{-d}{dr} [\frac{1}{R} \exp(-\mu r)] = Yukawa \text{ force}$$

r = radial distance, $2\pi R =$ circumference of great circle.

NEW: Solution form-invariant if one goes to rotating frame Only relative motion of quasars and gyroscope spins is relevant.

Bottom Line: Exact Dragging of Inertial Axes Mach's Principle holds for Vorticity, strong or weak

Inertial axes exactly follow the weighted average of cosmic energy currents (in all reference frames):

$$\vec{\Omega}_{\rm gyro} = \langle \vec{\Omega}_{\rm matter} \rangle \equiv \int_0^\infty dr \ \vec{\Omega}_{\rm matter}(r) \ W(r),$$

$$W(r) = rac{1}{3} \, 16 \pi G_{
m N}(
ho +
ho) \, R^3 \, Y_{\mu}(r).$$

This is the **exact solution** for vorticity on **Uniform-Radial-Expansion Slices** which are **spatially open**.

► Resulting weight function W(r) in exact solution normalized to unity ⇔ exact rotational frame-dragging.