

A new crystalline phase in magnetar crusts

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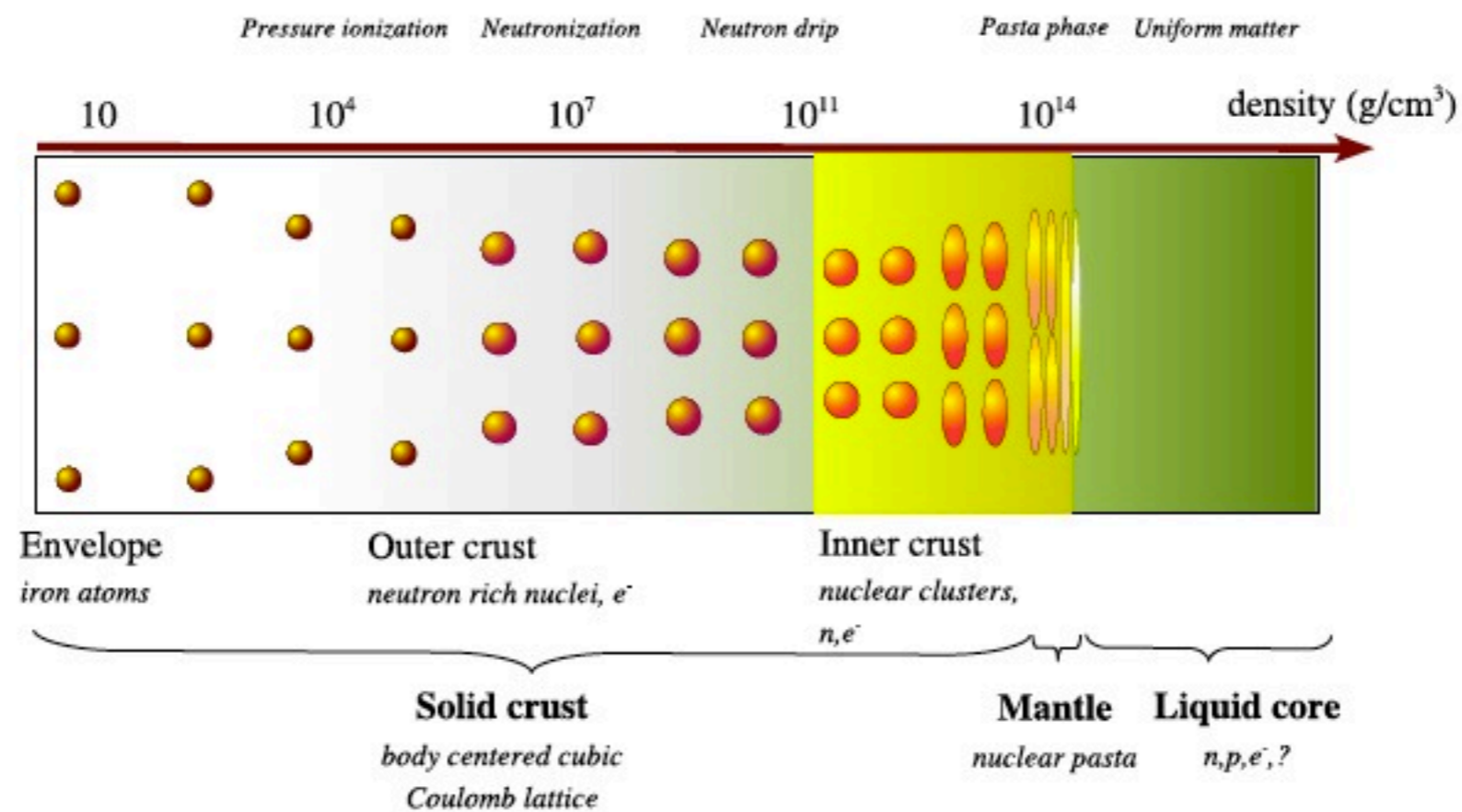
University of Maryland

Dec. 2013

Bedaque, Mahmoodifar, Sen, PRC 88, 055801 (2013)

Bedaque, Mahmoodifar, Ng, Sen [arXiv:1312.0591 (astro-ph.HE)]

Schematic picture of the ground state structure of neutron stars



- Large magnetic fields can change the structure and properties of matter
- Magnetars are highly magnetized neutron stars with $B \sim 10^{15}$ G
- Electrons motion perpendicular to the magnetic field will be quantized into Landau orbitals

Single electron energy levels:

$$E_n = \sqrt{k_z^2 + 2neB + m_e^2}$$

$$n_e = (eB/2\pi^2)k_e$$

We consider low enough densities where only the lowest Landau level is occupied.

$$\rho \lesssim \frac{AM}{\sqrt{2}\pi^2 Z} (eB)^{3/2} \approx 5.2 \times 10^8 \text{ g/cm}^3 \left(\frac{A_{66}}{Z_{28}} \right) B_{15}^{3/2}$$

Screening of the ion-ion potential by electrons

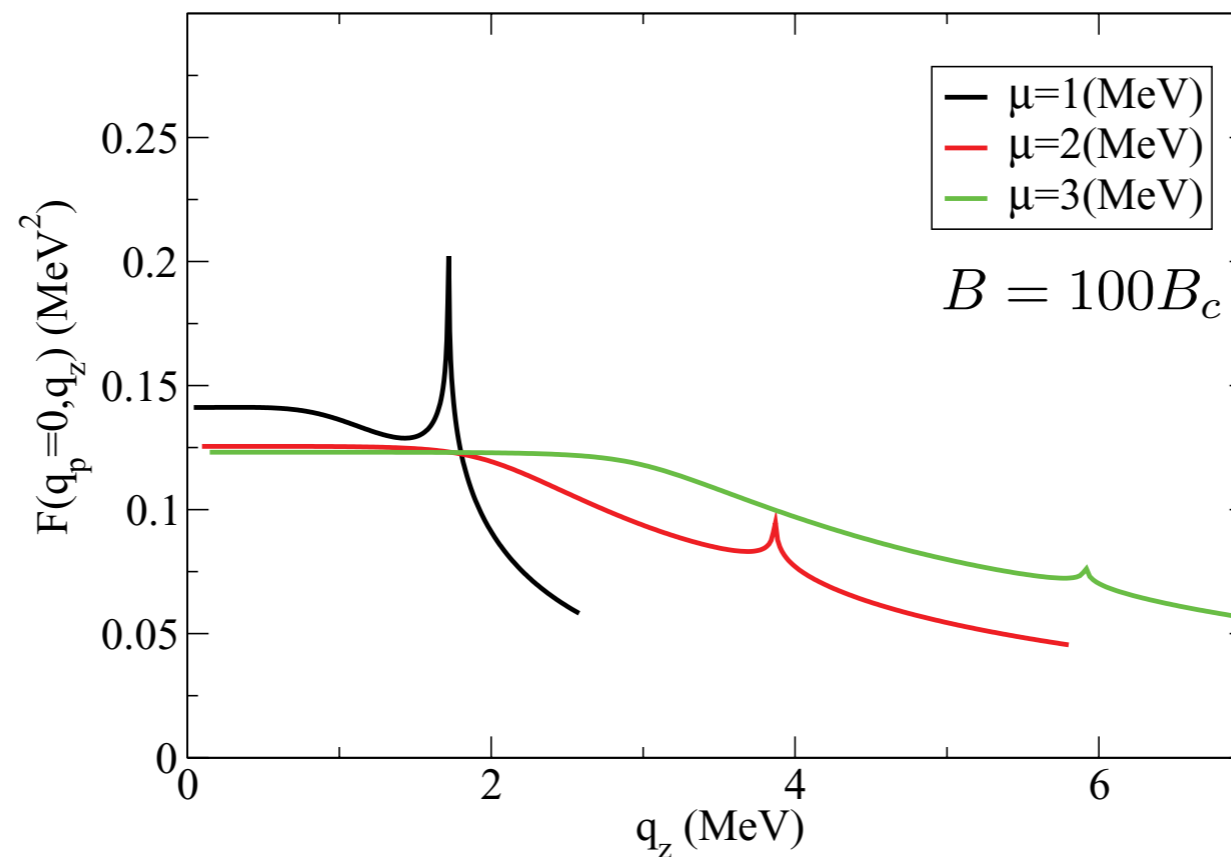
Sharma and Reddy (2011) computed the effect of one loop electron-hole polarization function on the ion-ion potential.

$$V(\mathbf{q}) = \frac{Z_1 Z_2 e^2}{q^2 - e^2 \Pi(\mathbf{q})}$$

$$V(\mathbf{r}) = \frac{Z_1 Z_2 e^2}{r} g(\mathbf{r})$$

$$g(\mathbf{r}) = 4\pi r \int \frac{d^3 q}{(2\pi)^3} \frac{e^{i\mathbf{q}\cdot\mathbf{r}}}{q^2 + F(\mathbf{q})}$$

Due to a sharp Fermi surface, $\Pi(q)$ shows non-analytic behavior at $q_z = \pm 2k_e$, which causes a long-range potential in the position space.

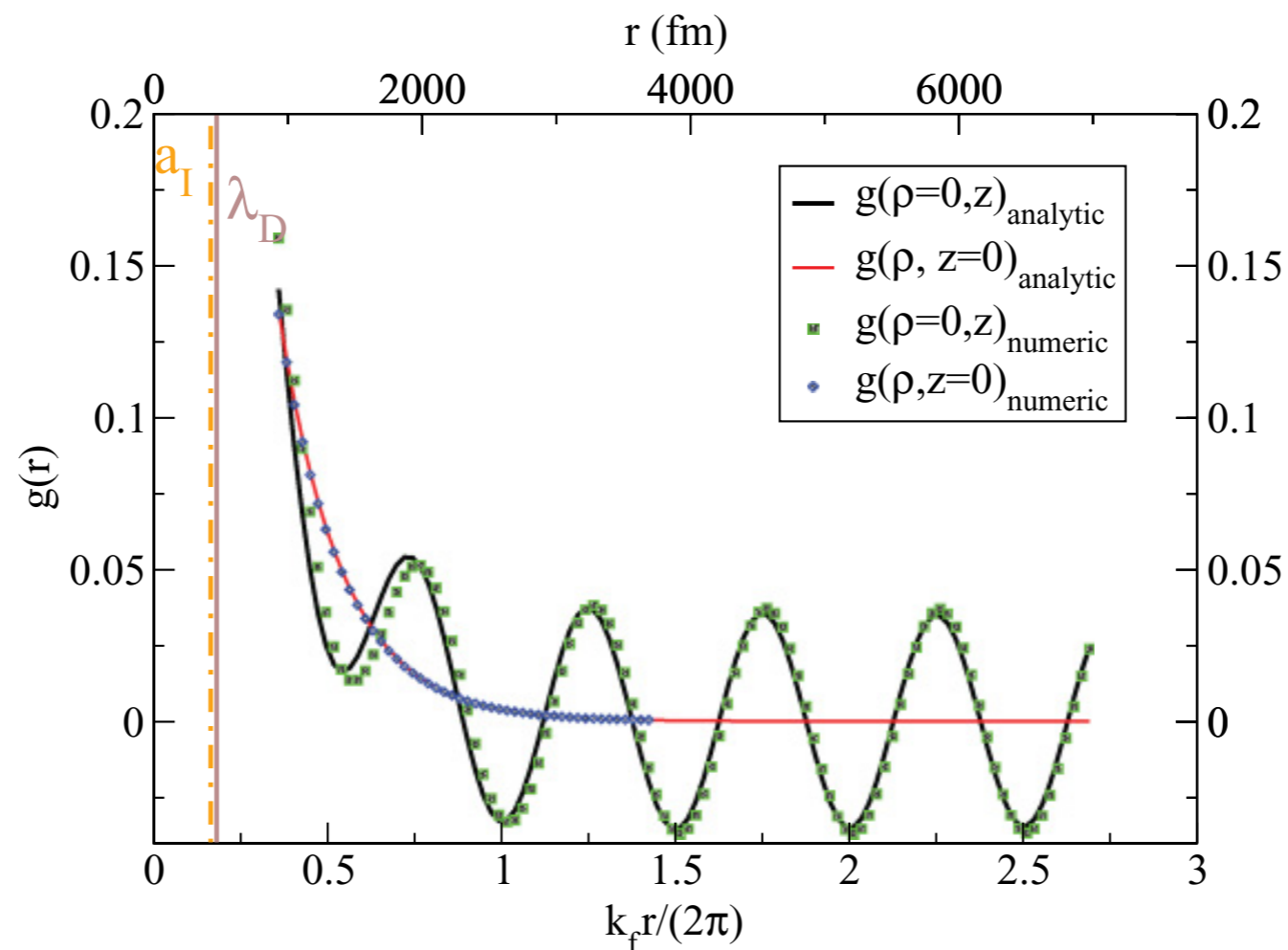


Sharma & Reddy, PRC **83**, 025803 (2011)

Only important when electrons are non-relativistic.

Friedel Oscillations

$$V(r_{\perp}, z) = Z^2 \alpha \left[\frac{e^{-m_D \sqrt{r_{\perp}^2 + z^2}}}{\sqrt{r_{\perp}^2 + z^2}} - \frac{m_D^2 e^{-z/\lambda_T}}{4z} \frac{\cos(2k_e z) r_{\perp}}{\sqrt{4k_e^2 + \frac{m_D^2}{2} \ln(4k_e z)}} K_1 \left(r_{\perp} \sqrt{4k_e^2 + \frac{m_D^2}{2} \ln(4k_e z)} \right) \right]$$



Sharma & Reddy, PRC **83**, 025803 (2011)

$$V_F(0, z) \approx -\frac{Z^2 \alpha m_D^2}{4} \frac{e^{-z/\lambda_T} \cos(2k_e z)}{z} \frac{1}{4k_e^2 + \frac{m_D^2}{2} \log(4k_e z)} f\left(\frac{2k_e^2 + \frac{m_D^2}{4} \log(4k_e z)}{eB}\right)$$

$$f(x) = 1 + x e^x E_i(-x),$$

$$V(\rho = 0, z) = Z^2 \alpha \left[\frac{e^{-m_D z}}{z} - \frac{m_D^2 e^{-z/\lambda_T}}{4z} \frac{\cos(2k_e z)}{4k_e^2 + \frac{m_D^2}{2} \ln(4k_e z)} \right]$$

$$\lambda_T = \frac{2\pi k_e}{mT} \approx (1.6 \times 10^{-7} \text{ cm}) \left(\frac{Z_{28}}{A_{66}}\right) \left(\frac{\rho_8}{B_{15} T_1}\right),$$

$$\frac{1}{m_D} \approx (1.3 \times 10^{-10} \text{ cm}) \sqrt{\frac{Z_{28} \rho_8}{A_{66}}} \frac{1}{B_{15}},$$

$$\frac{1}{k_e} \approx (3.0 \times 10^{-11} \text{ cm}) \left(\frac{A_{66}}{Z_{28}}\right) \left(\frac{B_{15}}{\rho_8}\right),$$

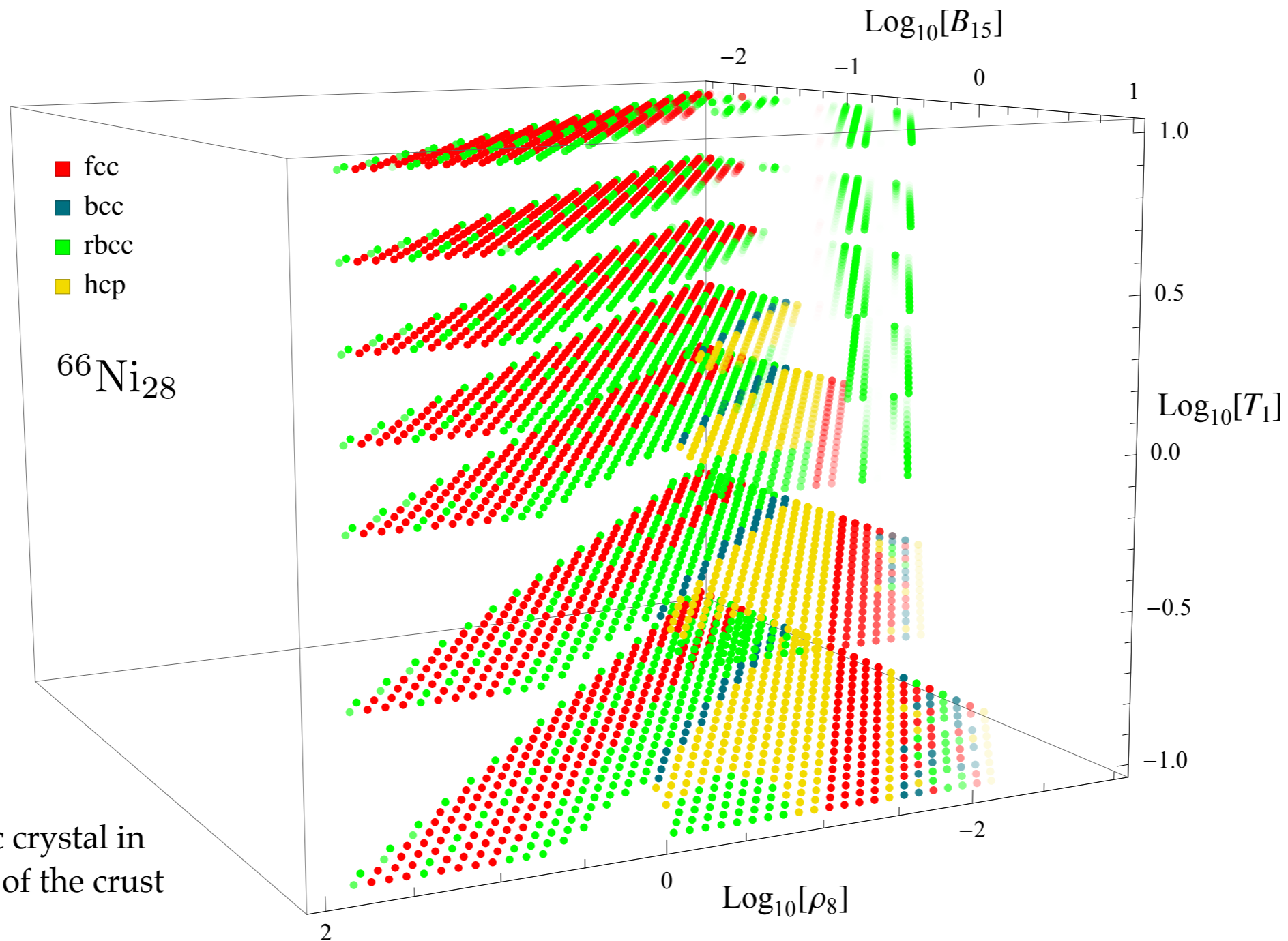
$$\frac{1}{\sqrt{eB}} \approx (8.1 \times 10^{-12} \text{ cm}) \frac{1}{\sqrt{B_{15}}}.$$

Bedaque, Mahmoodifar, Sen, PRC 88, 055801 (2013)

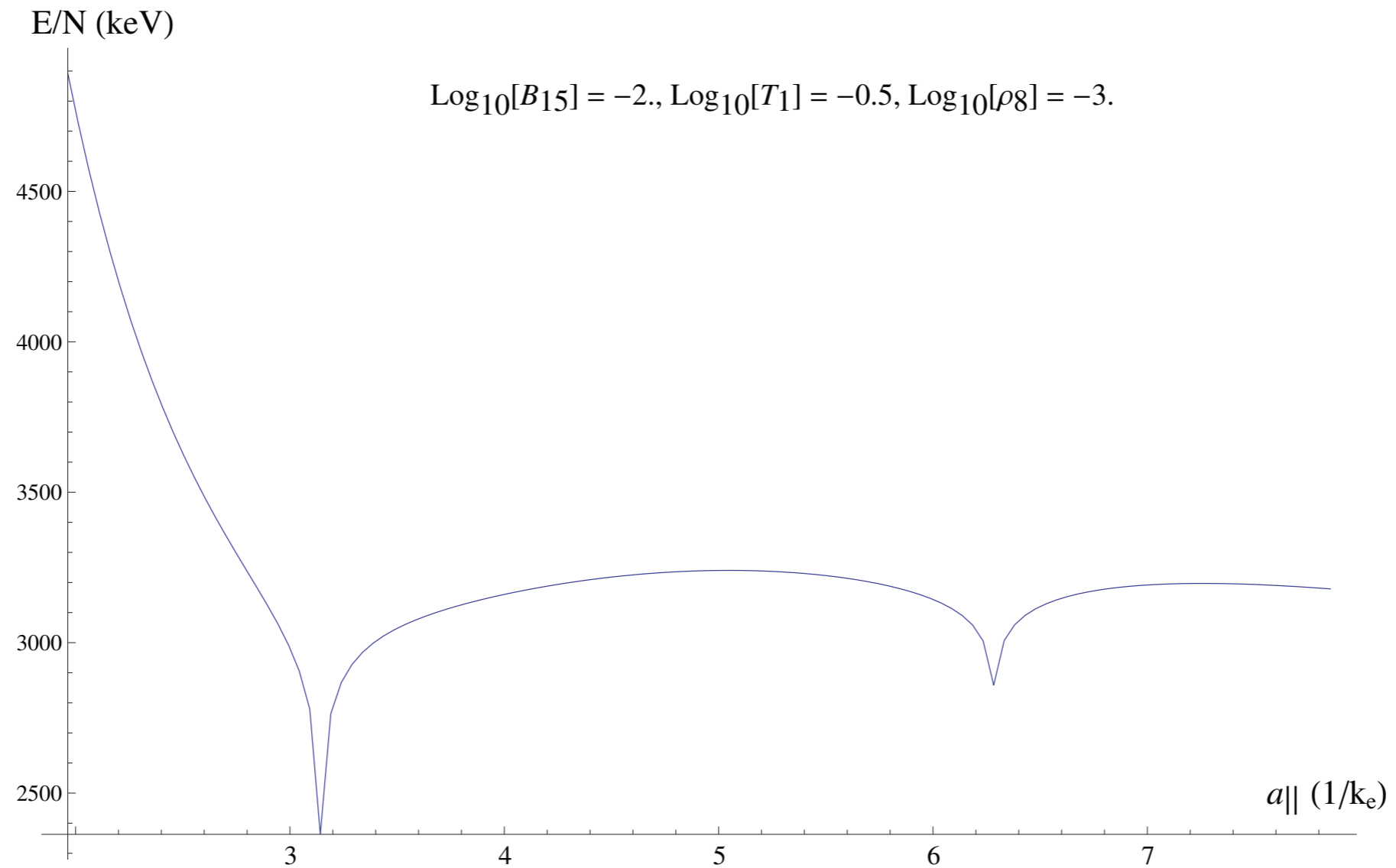
Parameter Space of “Friedel Crystals”

$$a_{\parallel} \sim n \times \frac{\pi}{k_e}$$

To keep the density fixed, a_{\perp} varies with a_{\parallel} .



Finite temperature effects:



$$\frac{E}{N} = \frac{1}{2} \sum_{i \neq 0} V(r_{\perp i}, z_i).$$

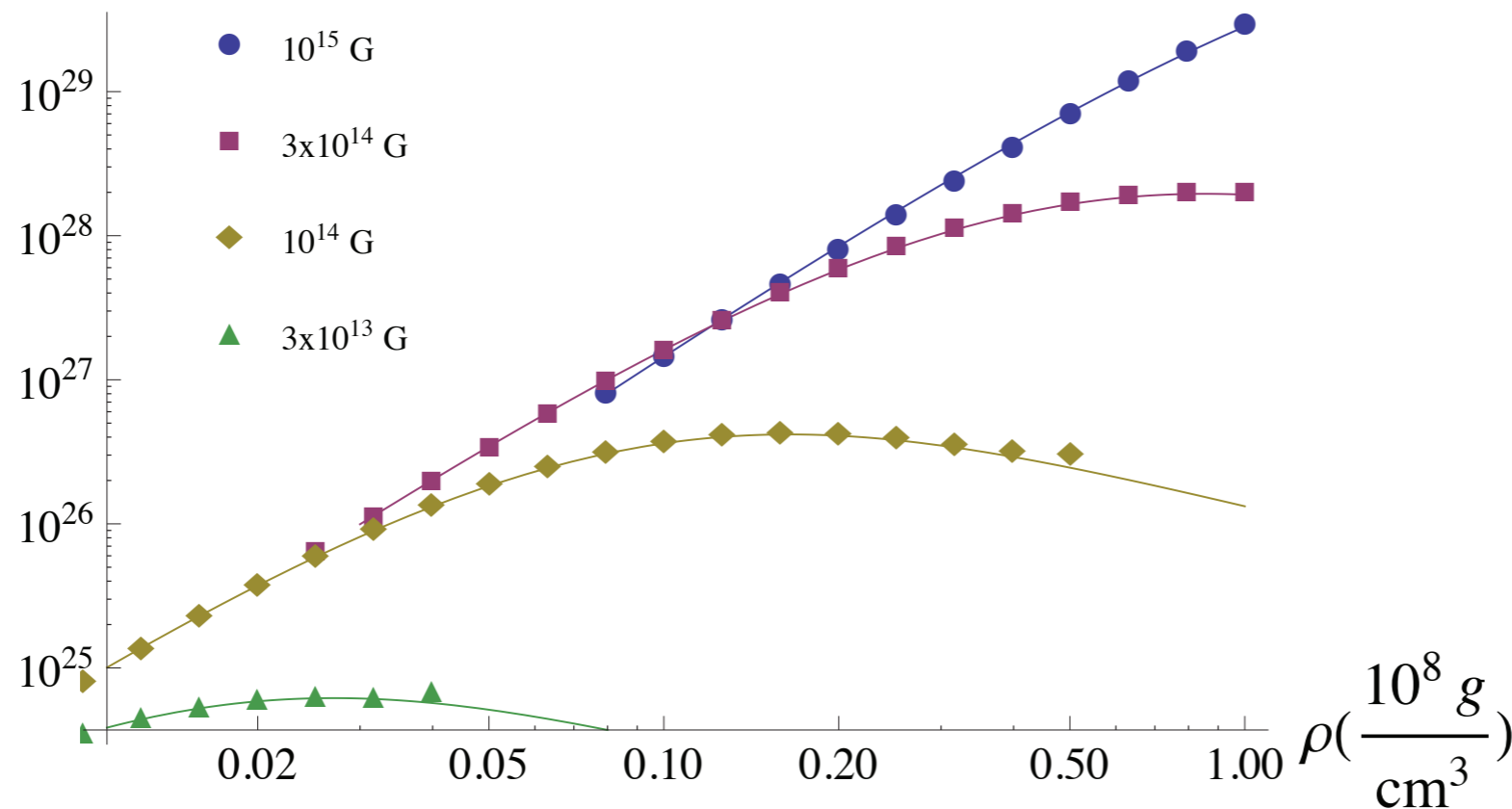
$$\lambda_{zz}^{xx}$$

The energy change of the lattice due to a small displacement in the x direction when the wave is propagating in the z direction.

$$\Delta U \approx \frac{1}{2} \int d^3x \lambda_{kl}^{ij} \partial_k \xi^i \partial_l \xi^j$$

$$\lambda_{zz}^{xx} \left(\frac{\text{ergs}}{\text{cm}^3} \right)$$

$$\delta F = \frac{1}{2} (c_{11} - c_{12}) u_{ii}^2 + c_{44} u_{ik} u_{ik} \quad (i \neq k)$$

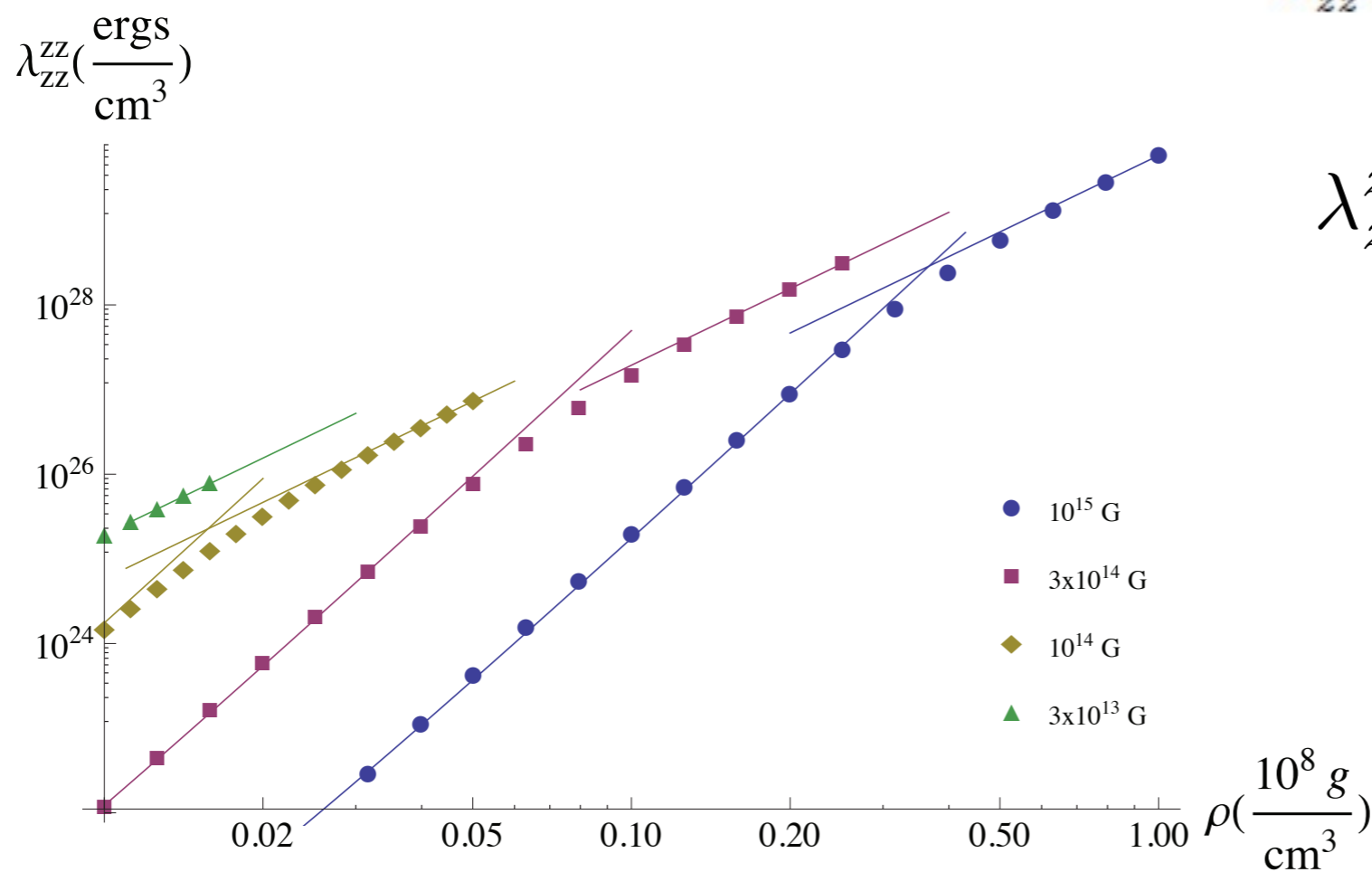


$T \sim 10^7 \text{ MeV}$

${}^{66}\text{Ni}_{28}$

λ_{zz}^{zz}

The energy change of the lattice due to a small displacement in the z direction when the wave is also propagating in the z direction.

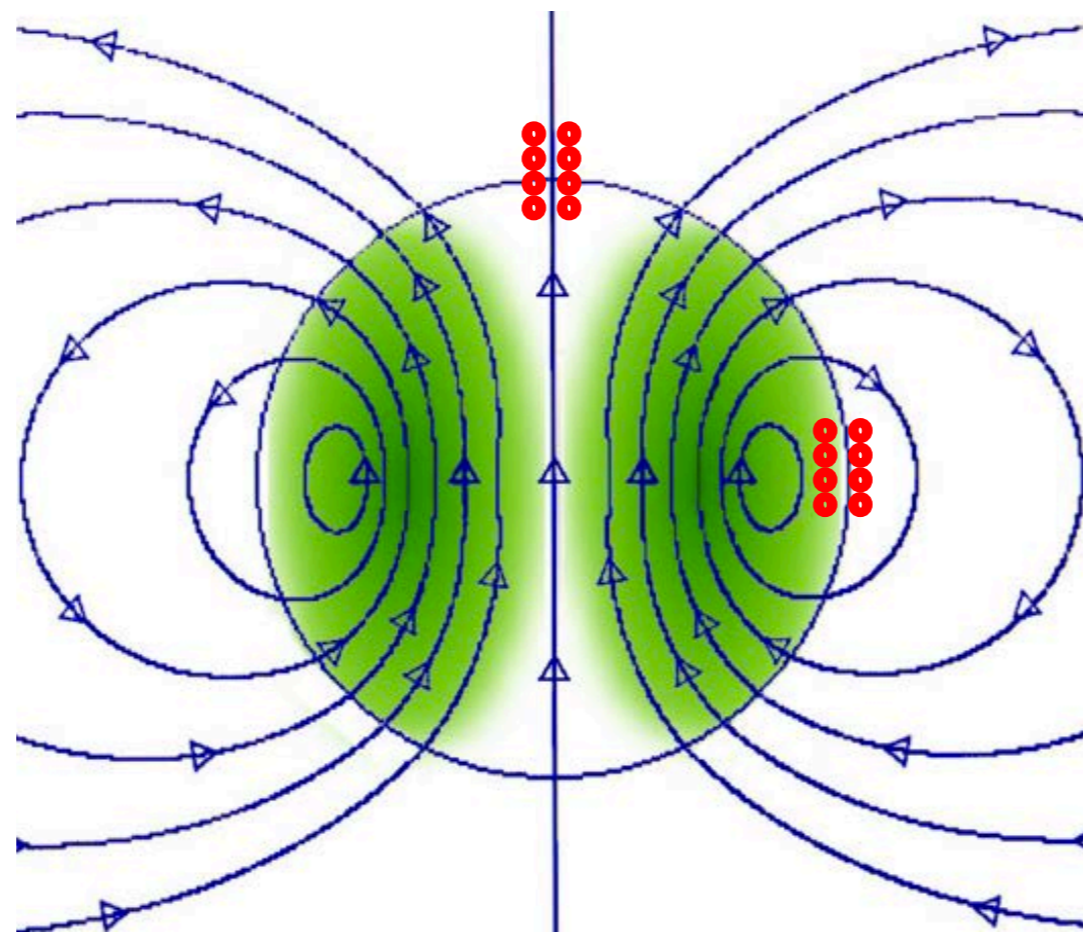


$$\lambda_{zz}^{zz} \approx 5.9 \times 10^{29} \frac{\text{ergs}}{\text{cm}^3} \left(\frac{Z_{28}^4 \rho_8^3}{A_{66}^3 B_{15} T_1^2} \right)$$

$$\lambda_{zz}^{zz} \sim 10^4 \times (c_{11} - c_{12})$$

$T \sim 10^7 \text{ MeV}$ ${}^{66}\text{Ni}_{28}$

Elastic constants that are dominated by the longitudinal structure of the lattice are significantly larger than that of a bcc Coulomb crystal of comparable densities.

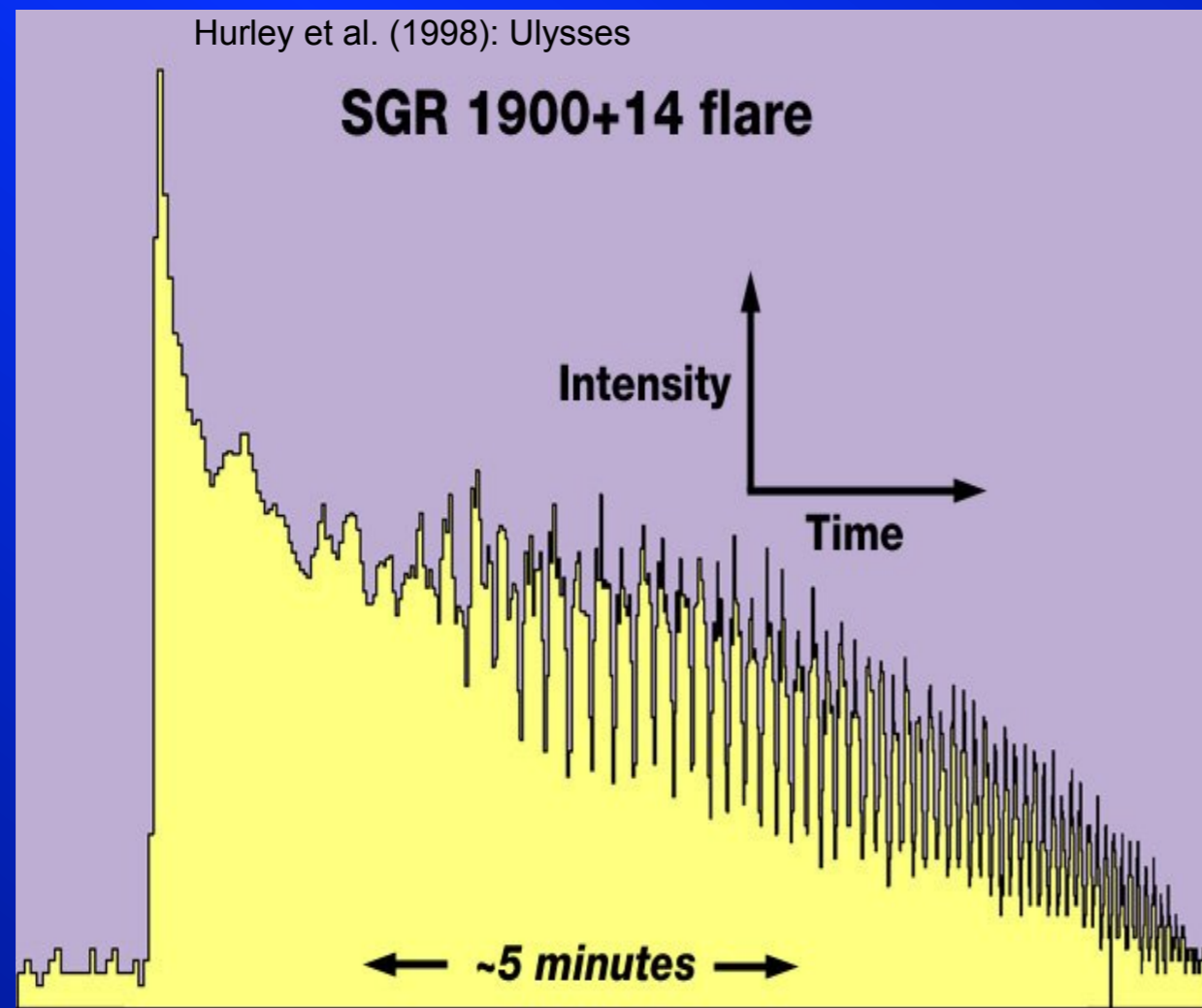


Potentially interesting implications for the X-ray oscillations seen from magnetars during their giant flares.



Goddard Space
Flight Center

Anatomy of a Hyperflare



Three events to date:

- March 5th 1979: SGR 0526-66
- August 27th 1998: SGR 1900+14
- December 27th 2004: SGR 1806-20

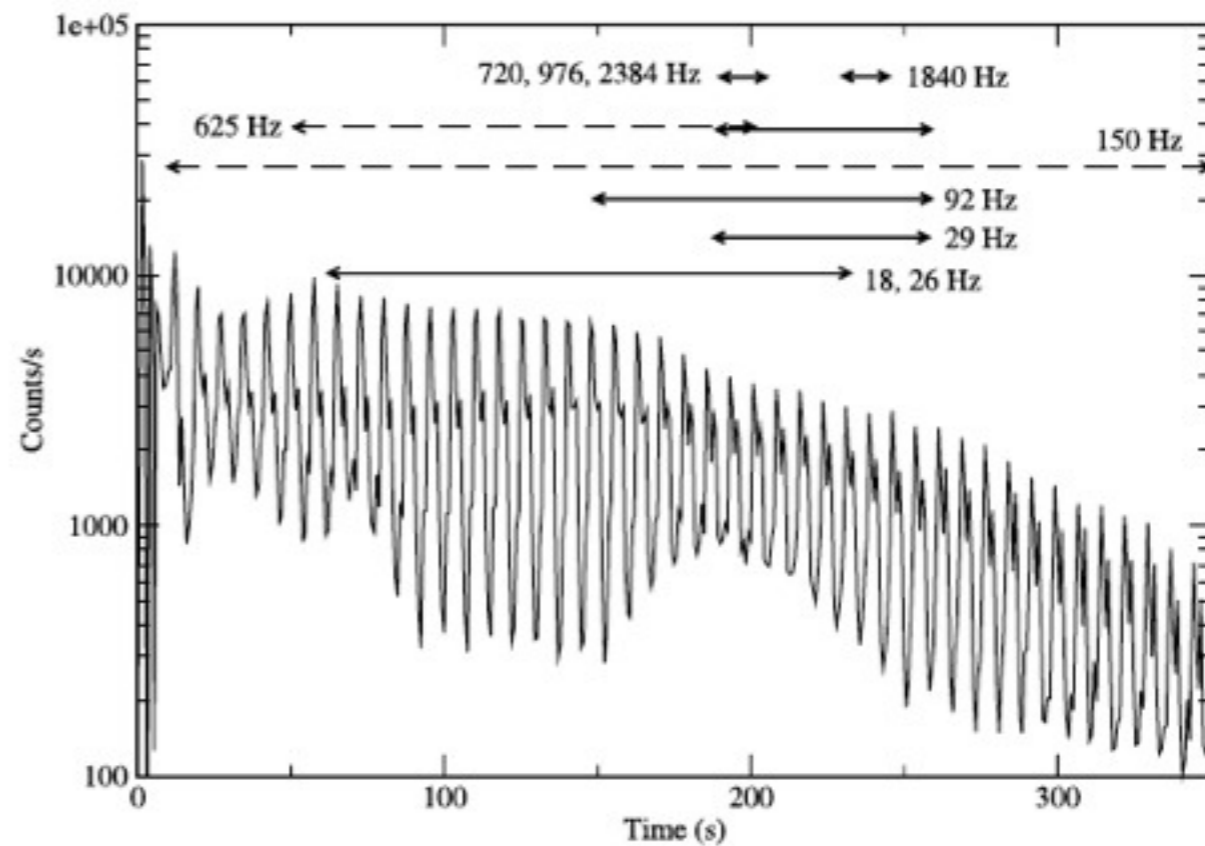
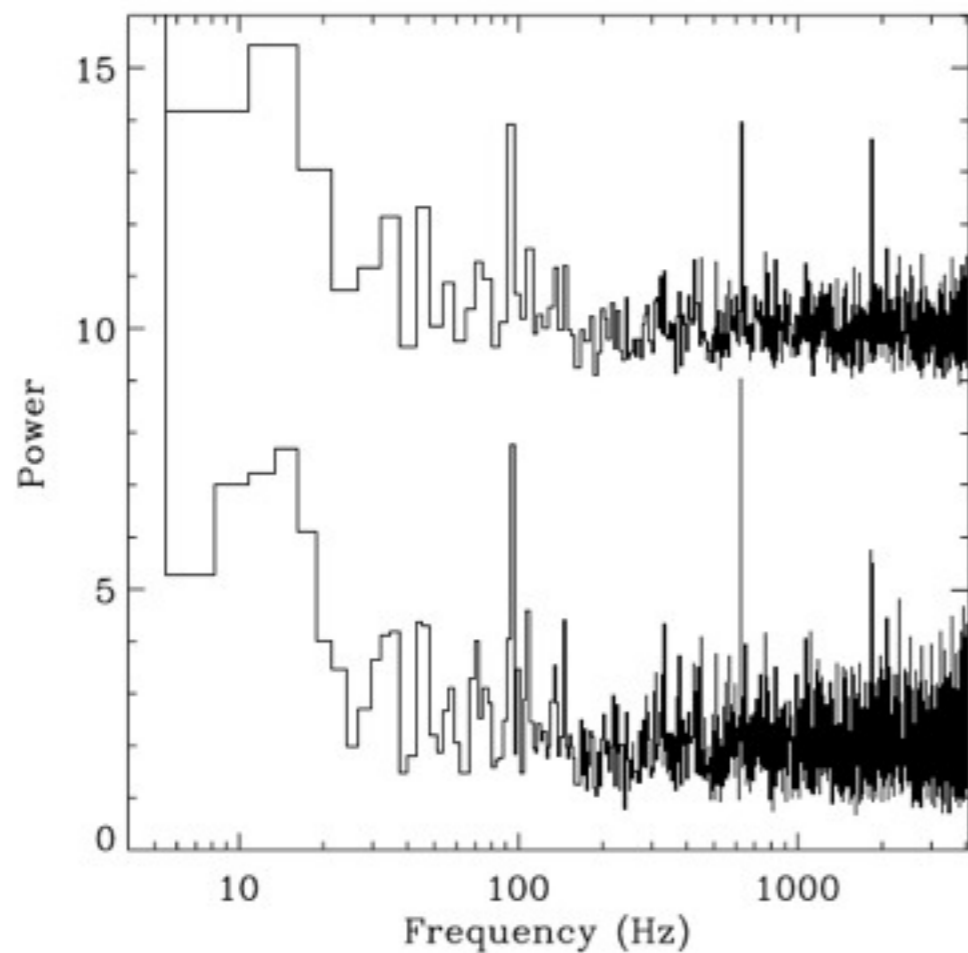
Powered by global magnetic instability (reconfiguration), crust fracturing.

- 10^{15} G magnetic fields implied (Thompson & Duncan 95)

- Short, hard, luminous initial pulse.
- Softer X-ray tail persists for minutes, and reveals neutron star spin period.
- Emission from a magnetically confined plasma.

Thanks to Tod strohmayr!

Dec. 2004 hyperflare from SGR 1806-20



T. Strohmayer and A. Watts 2006

The observations of **global oscillations** of neutron stars can provide a powerful **probe of their interior properties**, similar to the field of helioseismology.

SGR 1806–20		SGR 1900+14	
f (Hz)	Mode	f (Hz)	Mode
29	$0t_2$	28 ± 0.5	$0t_2$
92.7 ± 0.1	$0t_6$	53.5 ± 0.5	$0t_4$
150.3	$0t_{10}$	84	$0t_6$
626.46 ± 0.02	$1t_1$	155.1 ± 0.2	$0t_{11}$

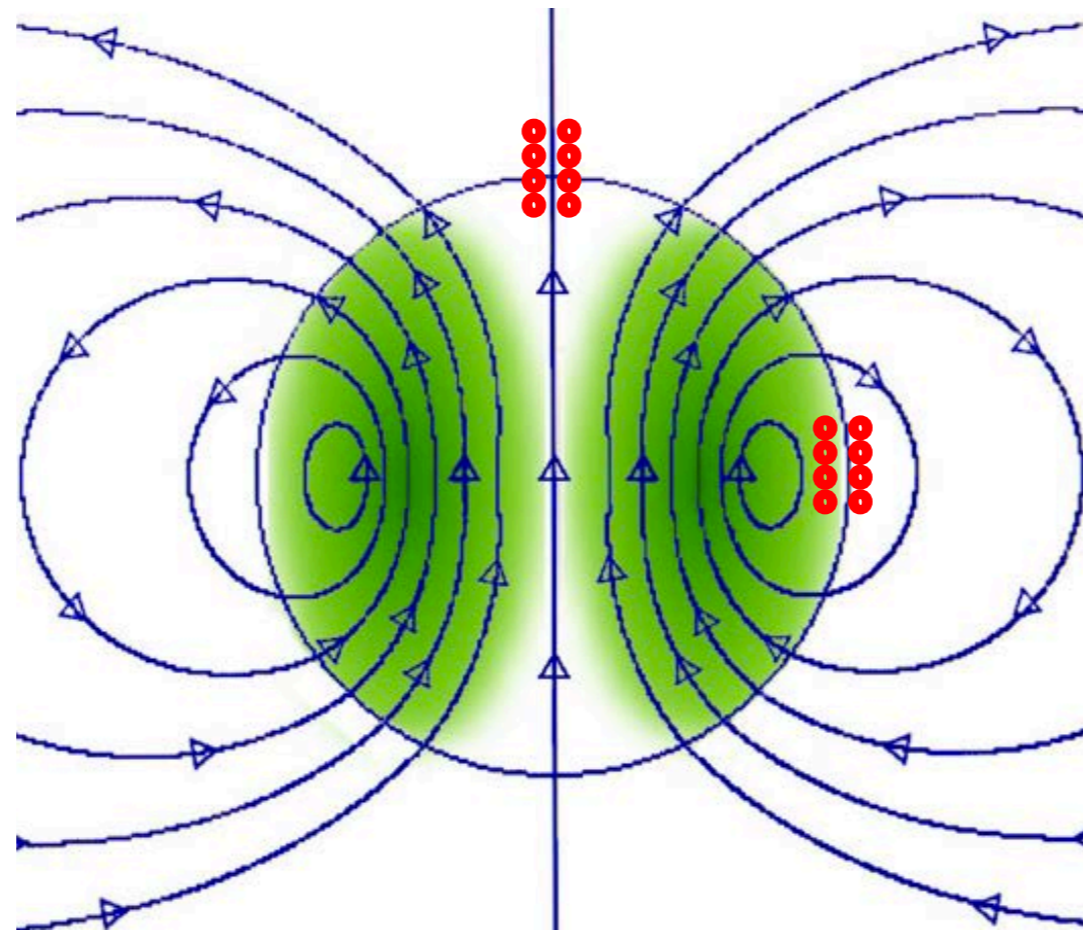
$$\omega_{0,l}^2 \propto \frac{v_t^2 l(l+1)}{R^2}$$

$$v_t = \sqrt{\mu/\rho}$$

Core Alfven modes?

Crust shear modes?

Effect of an anisotropic outer crust on the oscillation frequencies of the crustal modes?



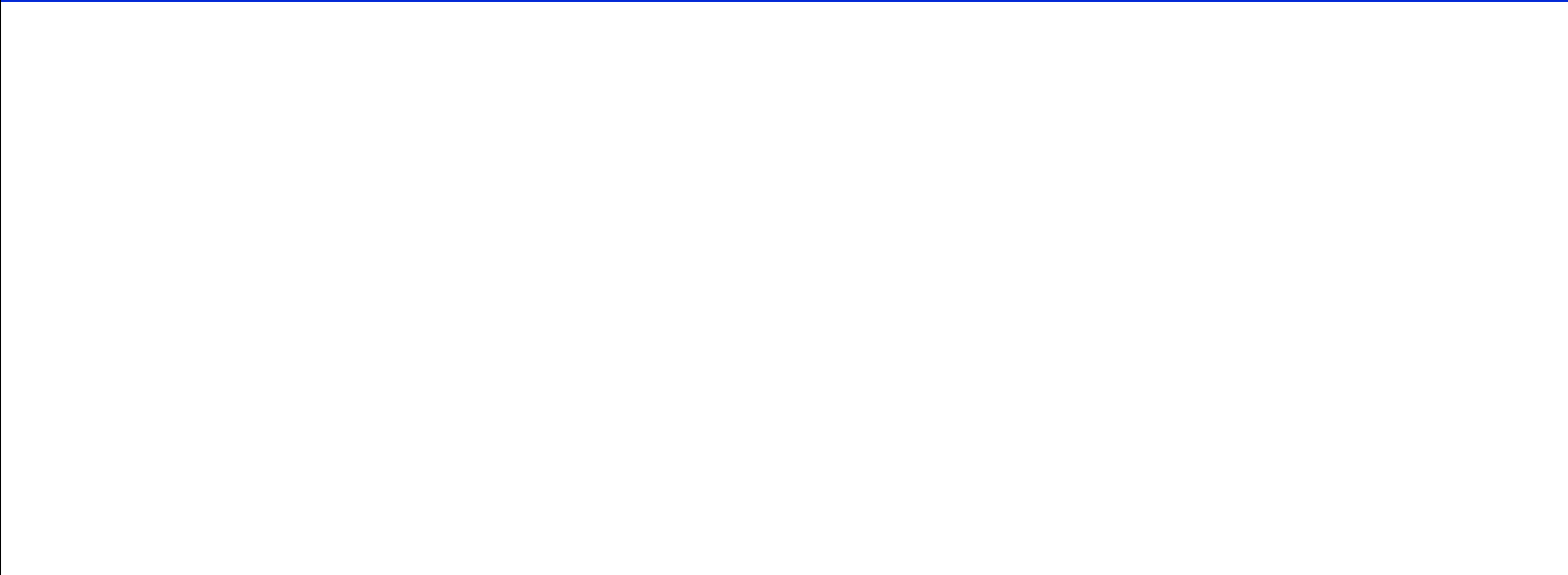
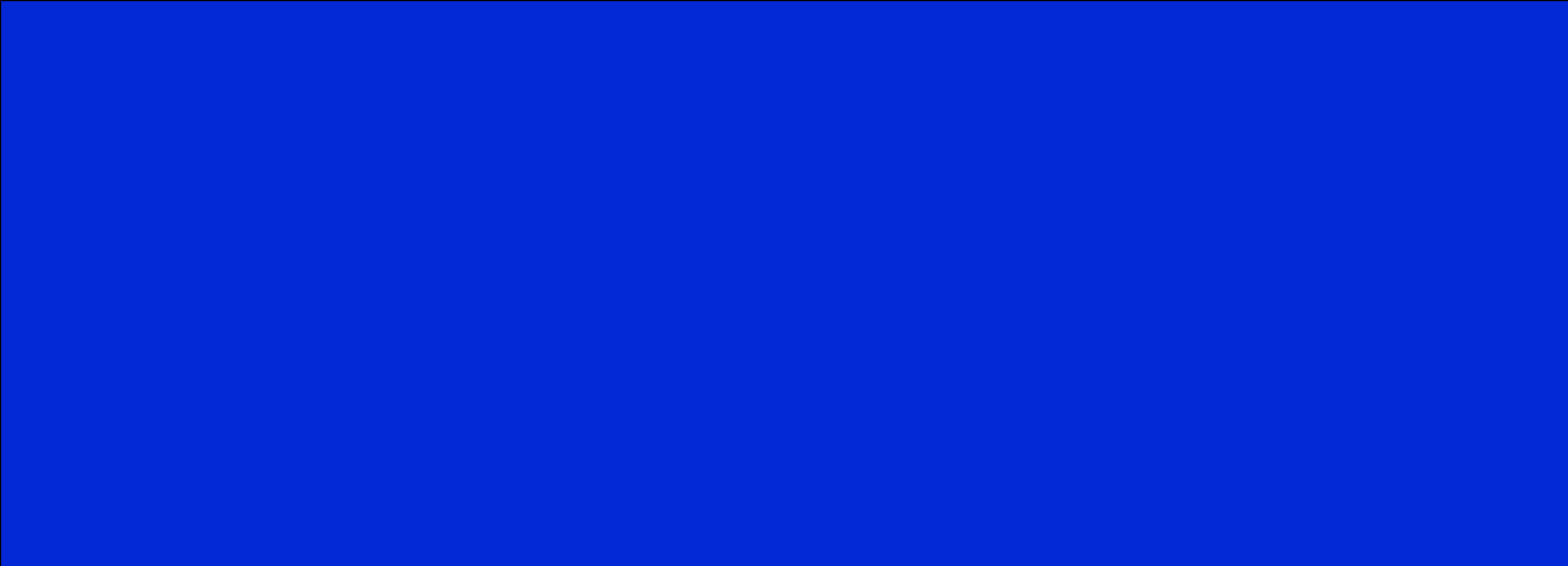
$$\mu_{eff} = 0.1194 \frac{n_i Z^2 e^2}{a}$$
$$\mu_{eff} = 0.1108 \frac{n_i Z^2 e^2}{a}$$

Ogata et al. 1990, Strohmayer et al.

Horowitz and Hughto

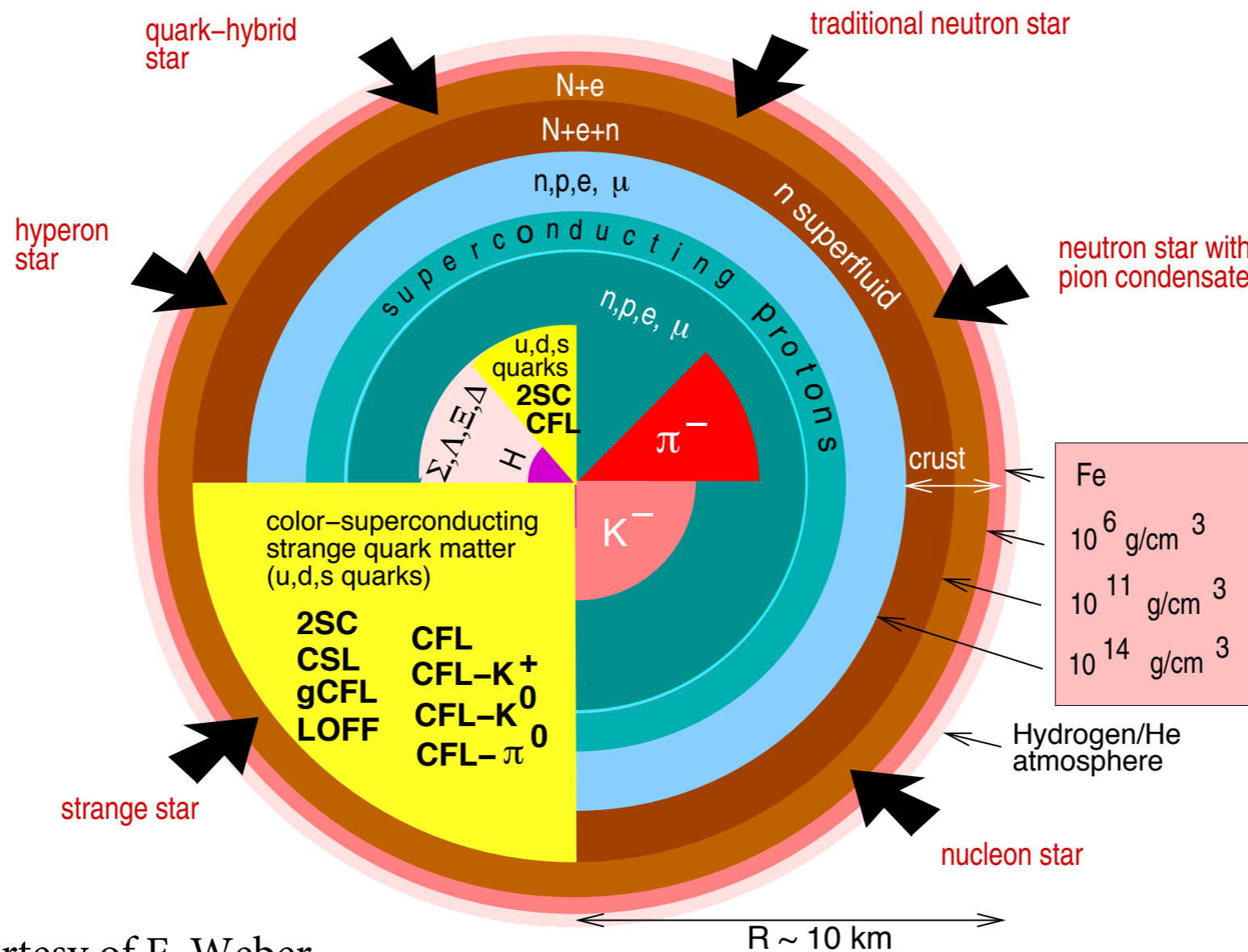
Conclusion

- Long range oscillations in the ion-ion potential along the magnetic field due to anisotropic screening of the Coulomb force by electrons in the presence of strong magnetic fields
- The long-ranged potential forces the ions to organize themselves into strongly coupled filaments along the magnetic field.
- Friedel crystals form in the outer crust of magnetars
- Large elastic constants in the longitudinal direction
- Implications for QPOs (shear mode frequencies) and GWs (breaking strain)

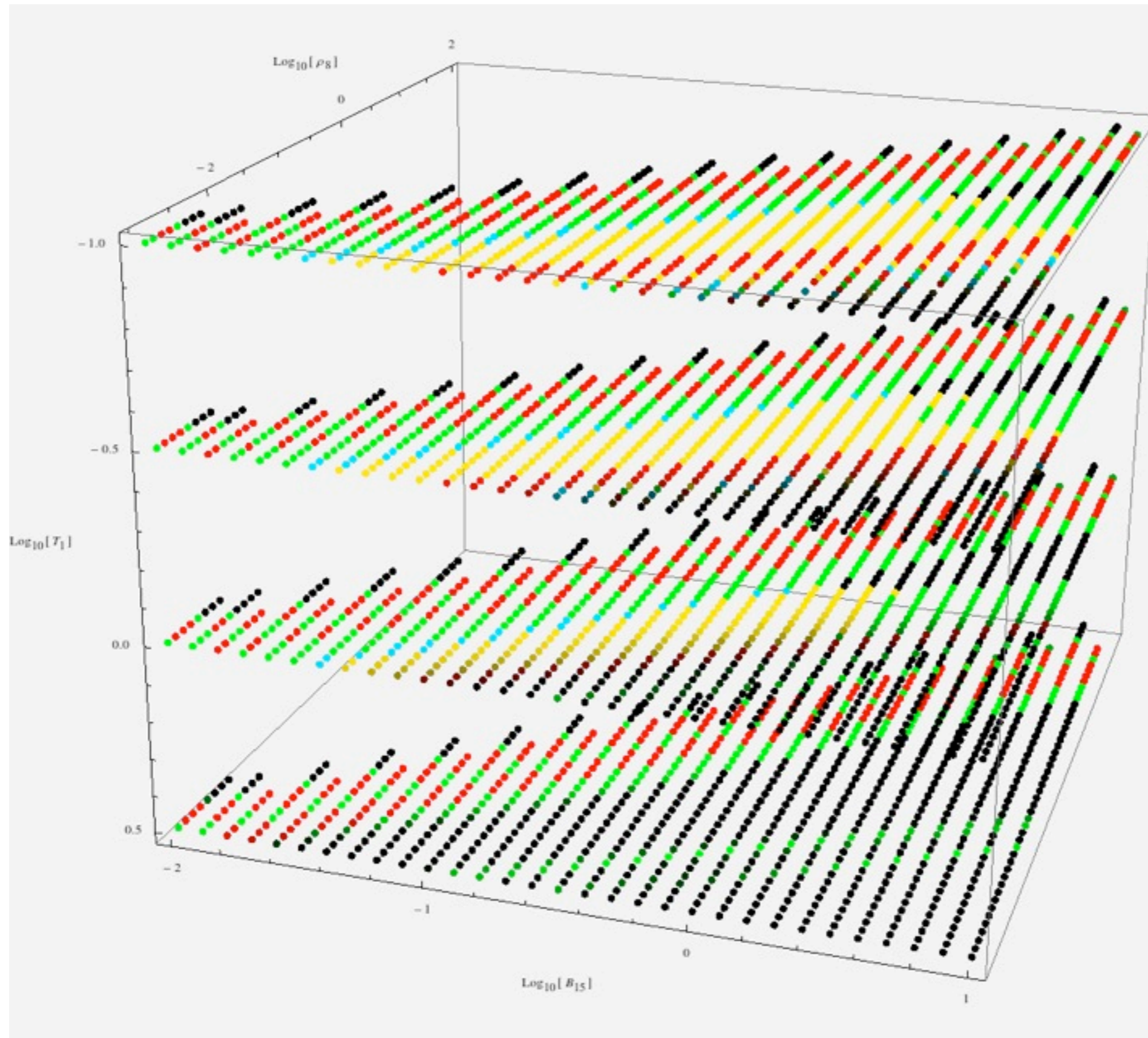


Compact Stars: The only “laboratory” for the study of cold ultra-dense matter

They may contain exotic form of matter.



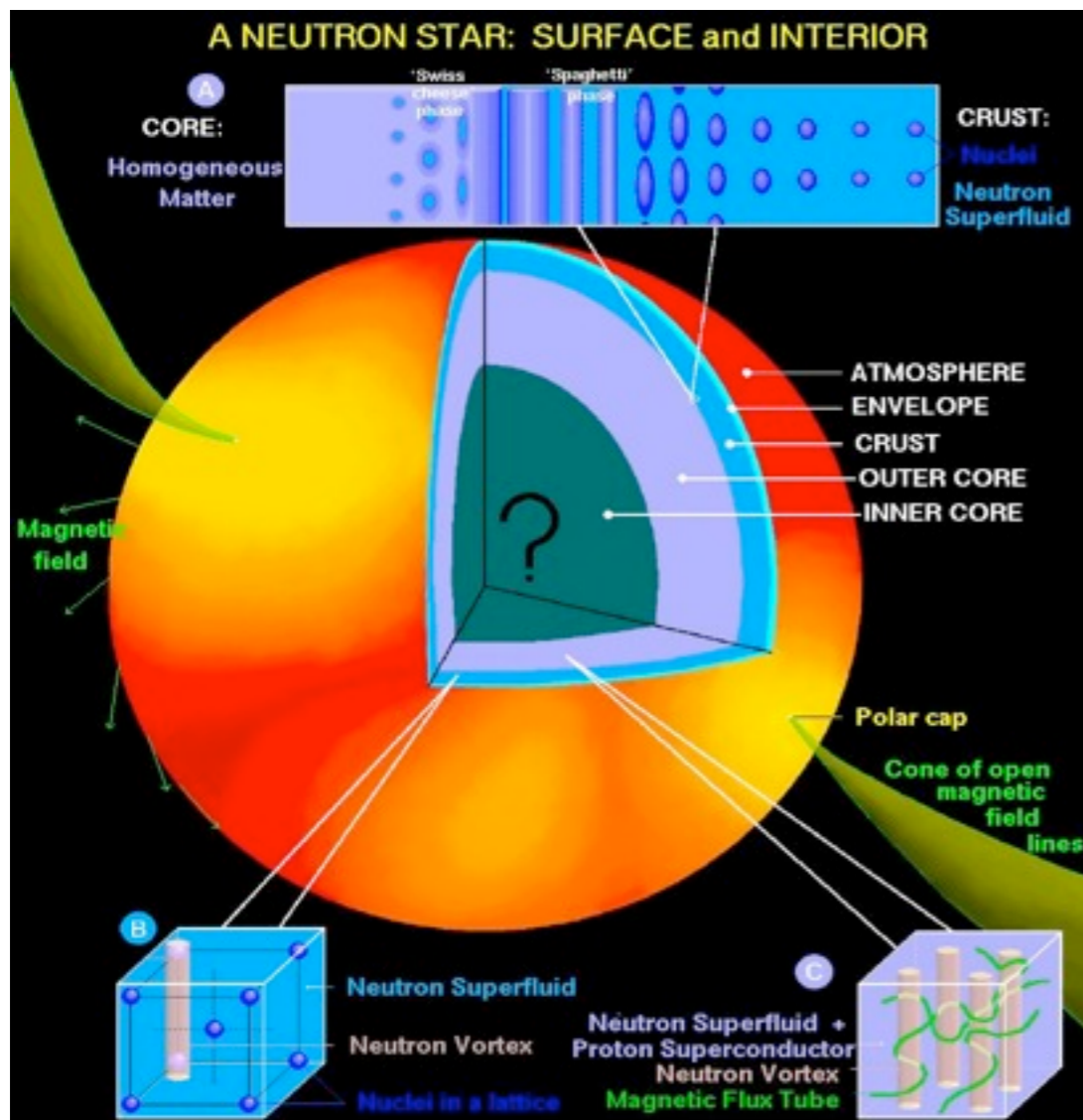
Courtesy of F. Weber



Nathan Ng

Compact Stars

The only “laboratory” for the study of cold ultra-dense matter



Mass $\sim 1.4M_{\odot}$

Radius $\sim O(10 \text{ km})$

Density $> \rho_{nuclear}$

$T < 1 \text{ MeV}$

They may contain exotic form of matter.

<http://www.astroscu.unam.mx/neutrones/NS-Picture/NStar/NStar.html>