

# Testing modified gravity at large distances with rotation curves

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# The Grumiller's Theory

In Grumiller's modified gravity theory, the spacetime is described by the EH action with the spherically symmetric metric

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta + \Phi^2 (d\theta^2 + \sin^2 \theta d\phi^2), \quad (1)$$

where the two dimensional metric  $g_{\alpha\beta}(x^\gamma)$  and the surface radius  $\Phi(x^\gamma)$  depends only on the coordinates  $x^\gamma = \{t, r\}$ . The 4-dimensional spherical spacetime manifold  $M$  can be decomposed to  $M = L \otimes S$ , L: radial, S: angular. The "spherical reduction" process simplifies the 4-dimensional EH action to a specific 2-dimensional dilation gravity model, demanding the IR-theory to be power-counting renormalizable and non-singularity of curvature at large  $\Phi$ . One obtains

$$S = - \int d^2x \sqrt{-g} [f(\Phi)R + 2(\partial\Phi)^2 - 2V(\Phi)], \quad (2)$$

where

$$V(\Phi) \equiv 3\Lambda\Phi^2 - 4a\Phi - 1. \quad (3)$$

## The Grumiller's Theory

The solutions, back in 4-d, are:

$$g_{\mu\nu} = \text{diag}(-K^2, 1/K^2, r^2, r^2 \sin^2 \theta) \quad \text{with} \quad K^2 \equiv 1 - \frac{2M}{r} - \Lambda r^2 + 2ar. \quad (4)$$

For a point particle with energy  $E$  and angular momentum  $\ell$  moving along a geodesic in the plane  $\theta = \pi/2$  in the background of the metric Eq.(4) one has:

$$\frac{1}{2} \left( \frac{dr}{d\tau} \right)^2 = \frac{1}{2} E^2 - V^{\text{eff}}, \quad (5)$$

where

$$V^{\text{eff}} \equiv \frac{K^2}{2} \left( 1 + \frac{\ell^2}{r^2} \right) = -\frac{M}{r} + \frac{\ell^2}{2r^2} - \frac{M\ell^2}{r^3} - \frac{\Lambda r^2}{2} + ar \left( 1 + \frac{\ell^2}{r^2} \right) + \dots \quad (6)$$

is the effective potential.

# Modified Newtonian Gravity

In Grumiller's approach the effective potential of a point mass ( $M$ ) is:

$$\phi_i(r) = -G \frac{M}{r} + a r, \quad (7)$$

where  $a$  is a universal constant, the Rindler acceleration.

For a smooth matter distribution with spherical symmetry one has:

$$\phi(r) = -G \int \frac{\rho(r')}{|\vec{r} - \vec{r}'|} d^3r' + a |\vec{r}|, \quad (8)$$

where  $\rho(r)$  is the density profile at radius  $r$ . The corresponding rotation velocity yields

$$v^2(r) = v_N^2(r) + v_R^2(r), \quad (9)$$

"N" is the Newtonian contribution and the Rindler velocity is

$$v_R^2(r) \equiv a |\vec{r}| \quad \text{or} \quad v_{GR}^2(r) \equiv a |\vec{r}|^n \quad (10)$$

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# The components

Bringing together all contributions to the total rotation curve, with stars ( $\star$ ), gas (G), modified gravity *a la* Grumiller (GR),

$$v_T^2(r) = \Upsilon_\star v_\star^2 + v_G^2 + v_{GR}^2(r), \quad (11)$$

where  $\Upsilon_\star$  is the mass-to-light ratio.

Note that there is no dark matter in the model !!



# Stars

The Newtonian contribution from stars ( $\star$ ) and gas ( $G$ ) can be thought as given by a spherized disk.

The contribution from stars is then given the Freeman-disk:

$$\rho_{\star}(r) = \frac{M_d}{2\pi r_d^2} e^{-r/r_d}, \quad (12)$$

where  $M_d$  is the mass of the disk and  $r_d$  its radius. Thus, the rotation curve contribution from stars with a standard Newtonian dynamics, yields

$$v_{\star}^2(r) = \frac{GM_d}{2r_d} \left(\frac{r}{r_d}\right)^2 \left[ I_0\left(\frac{r}{2r_d}\right) K_0\left(\frac{r}{2r_d}\right) - I_1\left(\frac{r}{2r_d}\right) K_1\left(\frac{r}{2r_d}\right) \right]. \quad (13)$$

On the other hand, the gas contribution  $v_G^2$  is computed by integrating its surface brightness as in standard Newtonian lore.

# Stars' mass models $\Upsilon_{\star}$

**Kroupa** Studies of the stellar population in the Milky Way suggest that the Kroupa IMF produces low disk masses that we consider as the minimal limit for the stellar disk.

**diet-Salpeter** Here  $\Upsilon_{\star}$  is set to be a constant value based on the diet-Salpeter IMF, in which the stellar population synthesis model has proven to give a maximum stellar disk for a given photometric constraint.

**Free  $\Upsilon_{\star}$**  Here we ignore a priori any knowledge of the IMF and treat  $\Upsilon_{\star}$  as an extra free parameter in the model, and we let the program to choose the best-fitted value for  $\Upsilon_{\star}$ .

# The Sample

## THINGS DATA

Galaxy (1)	Distance (2)	$R_d$ (3)	DIET-SALPETER	KROUPA
			$\log_{10} M_{star}$ (4)	$\log_{10} M_{star}$ (5)
N 925	9.2	3.30	10.01	9.86
N 2366	3.4	1.76	8.41	8.26
N 2403	3.2	1.81	9.67	9.52
N 2841	14.1	4.22	11.04	10.88
N 2903	8.9	2.40	10.15	10.0
N 2976	3.6	0.91	9.25	9.10
N 3031	3.6	1.93	10.84	10.69
N 3198	13.8	3.06	10.45	10.30
N 3521	10.7	3.09	11.09	10.94
N 3621	6.6	2.61	10.29	10.14
N 4736	4.7	1.99	10.27	10.12
N 5055	10.1	3.68	11.09	10.94
N 6946	5.9	2.97	10.77	10.62
N 7331	14.7	2.41	11.22	11.07
N 7793	3.9	1.25	9.44	9.29
I 2574	4.0	2.56	9.02	8.87
D 154	4.3	0.72	7.42	7.27

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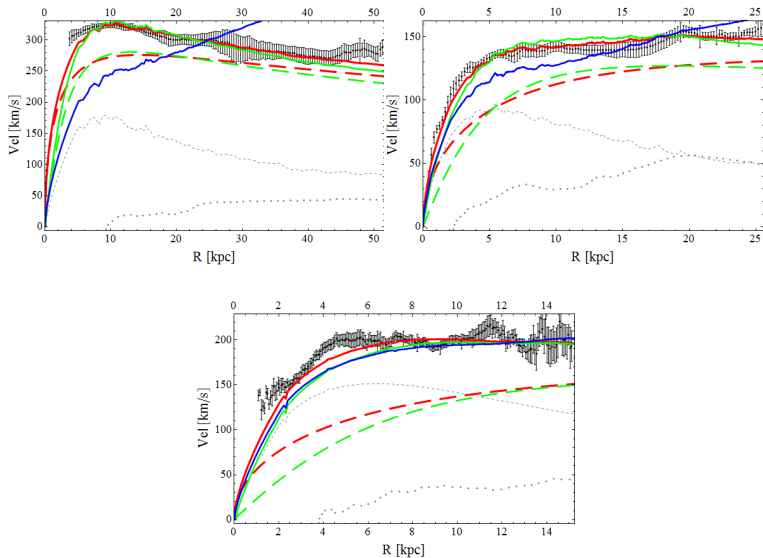
## Fitting Procedure

We make use of the observed rotation curve, stellar, and gas components as an input for the numerical code, in order to obtain the Rindler parameters  $(a, n)$ . To fit the observational velocity curve with the theoretical model we employ the  $\chi^2$  goodness-of-fit test ( $\chi^2$  test)

$$\chi^2 = \sum_{i=1}^n \left( \frac{v_{\text{obs}i} - v_{T_i}(r; a, n)}{\sigma_i} \right)^2, \quad (14)$$

where  $\sigma$  is the standard deviation, and  $n$  is the number of observations. One defines the reduced  $\chi_{\text{red}}^2 \equiv \chi^2 / (n - p - 1)$ , in which  $n$  is the number of observations and  $p$  is the number of fitted parameters.

# Results



## Results

Rindler  $n = 1$ 

Galaxy	Kroupa			diet-Salpeter		
	$a$	$a \left[ \frac{\text{cm}}{\text{s}^2} \right]$	$\chi_{\text{red}}^2$	$a$	$a \left[ \frac{\text{cm}}{\text{s}^2} \right]$	$\chi_{\text{red}}^2$
D 154	$358.45^{+5.45}_{-5.49}$	$1.16^{+0.02}_{-0.02}$	2.10	$355.64^{+5.45}_{-5.48}$	$1.15^{+0.02}_{-0.02}$	2.03
I 2574	$297.35^{+6.91}_{-6.97}$	$0.96^{+0.02}_{-0.02}$	4.74	$272.61^{+6.91}_{-6.97}$	$0.88^{+0.02}_{-0.02}$	5.53
N 2366	$285.45^{+1.94}_{-134.51}$	$0.93^{+0.01}_{-0.44}$	3.48	$285.45^{+2.66}_{-99.97}$	$0.93^{+0.01}_{-0.32}$	3.29
N 2403	$1258.65^{+5.77}_{-5.78}$	$4.08^{+0.02}_{-0.02}$	11.80	$1225.23^{+5.75}_{-5.76}$	$3.97^{+0.02}_{-0.02}$	9.29
N 2841	$2952.36^{+19.45}_{-19.50}$	$9.57^{+0.06}_{-0.06}$	76.10	$2405.18^{+18.91}_{-18.96}$	$7.79^{+0.06}_{-0.06}$	44.30
N 2903	$1985.81^{+13.91}_{-13.96}$	$6.44^{+0.05}_{-0.05}$	71.80	$1998.74^{+13.92}_{-13.97}$	$6.48^{+0.05}_{-0.05}$	74.10
N 2976	$1000.00^{+2.64}_{-723.12}$	$3.24^{+0.01}_{-2.34}$	8.05	$698.04^{+44.46}_{-44.80}$	$2.26^{+0.14}_{-0.15}$	10.70
N 3031	$2617.86^{+32.68}_{-32.77}$	$8.48^{+0.11}_{-0.11}$	9.54	$191.44^{+29.52}_{-29.62}$	$0.62^{+0.10}_{-0.10}$	22.80
N 3198	$632.49^{+5.85}_{-5.87}$	$2.05^{+0.02}_{-0.02}$	14.00	$567.72^{+5.78}_{-5.80}$	$1.84^{+0.02}_{-0.02}$	11.50
N 3521	$684.59^{+36.21}_{-36.59}$	$2.22^{+0.12}_{-0.12}$	6.14	$562.41^{+35.10}_{-35.50}$	$1.82^{+0.11}_{-0.12}$	7.54
N 3621	$889.79^{+6.98}_{-7.00}$	$2.88^{+0.02}_{-0.02}$	11.20	$777.46^{+6.86}_{-6.88}$	$2.52^{+0.02}_{-0.02}$	3.29
N 4736	$1005.43^{+27.96}_{-28.11}$	$3.26^{+0.09}_{-0.09}$	20.50	$420.13^{+27.07}_{-27.24}$	$1.36^{+0.09}_{-0.09}$	8.74
N 5055	$559.90^{+8.71}_{-8.74}$	$1.81^{+0.03}_{-0.03}$	5.09	$399.76^{+8.30}_{-8.34}$	$1.30^{+0.03}_{-0.03}$	15.80
N 6946	$1650.37^{+16.51}_{-16.55}$	$5.35^{+0.05}_{-0.05}$	1.29	$908.03^{+16.12}_{-16.17}$	$2.94^{+0.05}_{-0.05}$	4.42
N 7331	$1938.02^{+24.45}_{-24.55}$	$6.28^{+0.08}_{-0.08}$	8.65	$1505.92^{+23.90}_{-24.00}$	$4.88^{+0.08}_{-0.08}$	29.60
N 7793	$1720.65^{+1.57}_{-358.31}$	$5.58^{+0.01}_{-1.16}$	10.50	$1688.18^{+37.17}_{-15.34}$	$5.47^{+0.12}_{-0.05}$	5.52
N 925	$612.29^{+14.40}_{-14.39}$	$1.98^{+0.05}_{-0.05}$	3.74	$453.50^{+14.26}_{-14.36}$	$1.47^{+0.05}_{-0.05}$	5.72

## Results

Rindler  $n \neq 1$ 

Galaxy	Kroupa			diet-Salpeter		
	$a$	$n$	$\chi^2_{\text{red}}$	$a$	$n$	$\chi^2_{\text{red}}$
D 154	$429.25^{+6.52}_{-6.56}$	$0.87^{+0.010}_{-0.010}$	1.77	$420.27^{+6.43}_{-6.47}$	$0.88^{+0.010}_{-0.010}$	1.75
I 2574	$77.96^{+1.74}_{-1.75}$	$1.72^{+0.011}_{-0.011}$	0.98	$53.57^{+1.28}_{-1.29}$	$1.87^{+0.011}_{-0.011}$	1.32
N 2366	$340.76^{+16.15}_{-16.42}$	$1.02^{+0.029}_{-0.028}$	3.27	$285.45^{+18.13}_{-11.09}$	$1.13^{+0.100}_{-0.009}$	3.23
N 2403	$4667.30^{+20.97}_{-21.01}$	$0.48^{+0.002}_{-0.002}$	2.27	$4000.25^{+18.46}_{-18.50}$	$0.53^{+0.002}_{-0.002}$	2.55
N 2841	$59071.70^{+333.82}_{-334.49}$	$0.04^{+0.002}_{-0.002}$	0.90	$38957.50^{+262.73}_{-263.27}$	$0.13^{+0.002}_{-0.002}$	1.23
N 2903	$32294.30^{+211.39}_{-212.01}$	$0.03^{+0.002}_{-0.002}$	7.99	$33026.10^{+215.04}_{-215.67}$	$0.02^{+0.002}_{-0.002}$	8.34
N 2976	$796.89^{+25.89}_{-26.15}$	$2.14^{+0.045}_{-0.044}$	1.15	$386.08^{+336.71}_{-0.34}$	$3.43^{+2.196}_{-0.001}$	10.60
N 3031	$9840.04^{+116.91}_{-117.17}$	$0.36^{+0.006}_{-0.006}$	5.12	$2.12^{+0.12}_{-0.12}$	$3.34^{+0.024}_{-0.023}$	20.30
N 3198	$2830.13^{+25.46}_{-25.54}$	$0.54^{+0.003}_{-0.003}$	7.77	$1531.59^{+15.37}_{-15.42}$	$0.69^{+0.003}_{-0.003}$	9.54
N 3521	$195.76^{+10.31}_{-10.44}$	$1.42^{+0.017}_{-0.016}$	5.82	$48.00^{+2.85}_{-2.89}$	$1.83^{+0.019}_{-0.018}$	6.84
N 3621	$3728.60^{+28.09}_{-28.17}$	$0.50^{+0.003}_{-0.003}$	0.87	$1880.38^{+16.30}_{-16.34}$	$0.70^{+0.003}_{-0.003}$	1.24
N 4736	$8377.42^{+176.93}_{-177.79}$	$< 0.002$	5.76	$4216.82^{+176.24}_{-176.98}$	$< 0.003$	4.11
N 5055	$322.28^{+5.04}_{-5.06}$	$1.16^{+0.004}_{-0.004}$	4.96	$11.81^{+0.24}_{-0.24}$	$2.01^{+0.000}_{-0.005}$	13.50
N 6946	$2005.70^{+20.02}_{-20.07}$	$0.92^{+0.004}_{-0.004}$	1.25	$57.86^{+1.00}_{-1.00}$	$2.11^{+0.007}_{-0.007}$	2.47
N 7331	$876.99^{+11.11}_{-11.15}$	$1.27^{+0.004}_{-0.004}$	8.21	$80.47^{+1.26}_{-1.27}$	$1.99^{+0.005}_{-0.005}$	26.00
N 7793	$2850.27^{+35.74}_{-35.91}$	$0.69^{+0.009}_{-0.009}$	4.48	$1884.05^{+26.61}_{-26.74}$	$0.92^{+0.009}_{-0.009}$	5.31
N 925	$100.24^{+2.31}_{-2.32}$	$1.83^{+0.010}_{-0.010}$	1.35	$20.43^{+0.60}_{-0.60}$	$2.41^{+0.012}_{-0.012}$	2.30



## Results

Galaxy	Free mass model										
	$a$	$a$ [ $\frac{\text{cm}}{\text{s}^2}$ ]	$n = 1$	$M_B$	$\chi^2_{\text{red}}$	$a$	$n \neq 1$	$M_D$	$M_B$	$\chi^2_{\text{red}}$	
D 154	332.91 <sup>+5.42</sup> <sub>-5.50</sub>	1.08 <sup>+0.02</sup> <sub>-0.02</sub>	7.96 <sup>+7.06</sup> <sub>-7.07</sub>		1.78	378.75 <sup>+6.01</sup> <sub>-6.05</sub>	0.93 <sup>+0.010</sup> <sub>-0.010</sub>	7.80 <sup>+7.06</sup> <sub>-7.06</sub>		1.75	
I 2574	365.24 <sup>+6.89</sup> <sub>-6.95</sub>	1.18 <sup>+0.02</sup> <sub>-0.02</sub>	<1		3.94	136.85 <sup>+2.63</sup> <sub>-2.65</sub>	1.52 <sup>+0.009</sup> <sub>-0.009</sub>	8.26 <sup>+7.76</sup> <sub>-7.77</sub>		1.03	
N 2366	176.84 <sup>+16.41</sup> <sub>-16.69</sub>	0.57 <sup>+0.05</sup> <sub>-0.05</sub>	9.00 <sup>+7.12</sup> <sub>-8.69</sub>		2.30	<10	1.11 <sup>+1.00</sup> <sub>-1.00</sub>	9.31 <sup>+9.26</sup> <sub>-7.92</sub>		1.58	
N 2403	797.22 <sup>+97.65</sup> <sub>-0.32</sub>	2.58 <sup>+0.32</sup> <sub>-0.30</sub>	10.2 <sup>+9.83</sup> <sub>-6.36</sub>	6.77 <sup>+5.80</sup> <sub>-5.72</sub>	4.88	3070.1 <sup>+16.1</sup> <sub>-16.1</sub>	0.59 <sup>+0.002</sup> <sub>-0.002</sub>	9.85 <sup>+8.07</sup> <sub>-8.07</sub>	7.70 <sup>+6.65</sup> <sub>-6.65</sub>	0.75	
N 2841	1182.6 <sup>+16.57</sup> <sub>-16.73</sub>	3.83 <sup>+0.05</sup> <sub>-0.05</sub>	11.4 <sup>+9.08</sup> <sub>-9.08</sub>		1.08	61227 <sup>+358</sup> <sub>-359</sub>	0.02 <sup>+0.002</sup> <sub>-0.002</sub>	11.0 <sup>+9.08</sup> <sub>-9.08</sub>		0.14	
N 2903	965.84 <sup>+12.81</sup> <sub>-12.87</sub>	3.13 <sup>+0.04</sup> <sub>-0.04</sub>	10.8 <sup>+8.74</sup> <sub>-8.75</sub>	7.43 <sup>+6.39</sup> <sub>-6.47</sub>	2.47	4686.8 <sup>+52.1</sup> <sub>-52.3</sub>	0.53 <sup>+0.004</sup> <sub>-0.004</sub>	10.7 <sup>+8.74</sup> <sub>-8.75</sub>	7.70 <sup>+6.71</sup> <sub>-6.77</sub>	1.82	
N 2976	2646.5 <sup>+99.93</sup> <sub>-21.12</sub>	8.58 <sup>+0.32</sup> <sub>-0.07</sub>	5.86 <sup>+6.95</sup> <sub>-4.95</sub>		2.32	2130.7 <sup>+41.2</sup> <sub>-41.6</sub>	1.30 <sup>+0.030</sup> <sub>-0.030</sub>	8.38 <sup>+7.59</sup> <sub>-7.59</sub>		0.99	
N 3031	1602.5 <sup>+31.69</sup> <sub>-31.80</sub>	5.19 <sup>+0.10</sup> <sub>-0.10</sub>	10.8 <sup>+8.50</sup> <sub>-8.50</sub>	7.83 <sup>+6.79</sup> <sub>-6.78</sub>	5.93	34102 <sup>+233</sup> <sub>-228</sub>	<0.002	10.5 <sup>+8.50</sup> <sub>-8.51</sub>	5.86 <sup>+4.88</sup> <sub>-4.89</sub>	4.86	
N 3198	503.00 <sup>+5.72</sup> <sub>-5.74</sub>	1.63 <sup>+0.02</sup> <sub>-0.02</sub>	10.5 <sup>+8.62</sup> <sub>-8.62</sub>	7.96 <sup>+6.95</sup> <sub>-6.97</sub>	5.11	2144.6 <sup>+21.0</sup> <sub>-21.0</sub>	0.60 <sup>+0.003</sup> <sub>-0.003</sub>	10.4 <sup>+8.62</sup> <sub>-8.62</sub>	8.65 <sup>+7.64</sup> <sub>-7.69</sub>	3.96	
N 3521	4580.0 <sup>+7.67</sup> <sub>-294</sub>	14.80 <sup>+0.02</sup> <sub>-0.95</sub>	9.00 <sup>+6.88</sup> <sub>-6.45</sub>		63.10	120.08 <sup>+6.75</sup> <sub>-8.96</sub>	1.51 <sup>+0.017</sup> <sub>-0.023</sub>	11.0 <sup>+8.96</sup> <sub>-9.05</sub>		1.78	
N 3621	689.06 <sup>+4.84</sup> <sub>-9.49</sub>	2.23 <sup>+0.02</sup> <sub>-0.03</sub>	10.4 <sup>+8.39</sup> <sub>-8.32</sub>		0.67	2015.0 <sup>+17.3</sup> <sub>-17.4</sub>	0.67 <sup>+0.003</sup> <sub>-0.003</sub>	10.3 <sup>+8.35</sup> <sub>-8.35</sub>		0.56	
N 4736	<10	<0.1	10.5 <sup>+8.44</sup> <sub>-8.44</sub>	7.27 <sup>+6.34</sup> <sub>-6.24</sub>	53.0	19022 <sup>+153</sup> <sub>-154</sub>	0.10 <sup>+0.001</sup> <sub>-0.001</sub>	7.08 <sup>+7.90</sup> <sub>-5.90</sub>	6.69 <sup>+5.77</sup> <sub>-5.63</sub>	32.1	
N 5055	556.69 <sup>+8.75</sup> <sub>-8.82</sub>	1.80 <sup>+0.03</sup> <sub>-0.03</sub>	11.0 <sup>+8.76</sup> <sub>-8.77</sub>	<1	9.30	27370 <sup>+175</sup> <sub>-175</sub>	0.02 <sup>+0.002</sup> <sub>-0.002</sub>	10.5 <sup>+8.76</sup> <sub>-8.76</sub>	7.70 <sup>+6.62</sup> <sub>-6.75</sub>	0.86	
N 6946	1378.7 <sup>+16.64</sup> <sub>-16.30</sub>	4.47 <sup>+0.05</sup> <sub>-0.05</sub>	10.7 <sup>+8.55</sup> <sub>-8.54</sub>	<1	1.48	7052.8 <sup>+53.5</sup> <sub>-53.6</sub>	0.50 <sup>+0.003</sup> <sub>-0.003</sub>	10.6 <sup>+8.54</sup> <sub>-8.54</sub>	3.70 <sup>+2.67</sup> <sub>-2.70</sub>	1.22	
N 7331	1707.9 <sup>+24.65</sup> <sub>-24.73</sub>	5.53 <sup>+0.08</sup> <sub>-0.08</sub>	11.0 <sup>+9.01</sup> <sub>-9.01</sub>	12.2 <sup>+11.2</sup> <sub>-11.1</sub>	0.26	3348.2 <sup>+45.7</sup> <sub>-45.8</sub>	0.79 <sup>+0.005</sup> <sub>-0.005</sub>	10.9 <sup>+9.01</sup> <sub>-9.01</sub>	7.70 <sup>+6.67</sup> <sub>-6.77</sub>	0.24	
N 7793	1523.1 <sup>+23.79</sup> <sub>-23.79</sub>	4.94 <sup>+0.08</sup> <sub>-0.08</sub>	9.47 <sup>+7.85</sup> <sub>-7.85</sub>		4.88	2277.8 <sup>+30.2</sup> <sub>-30.4</sub>	0.83 <sup>+0.009</sup> <sub>-0.009</sub>	9.31 <sup>+7.85</sup> <sub>-7.85</sub>		4.61	
N 925	803.09 <sup>+0.42</sup> <sub>-450</sub>	2.60 <sup>+0.01</sup> <sub>-1.46</sub>	7.55 <sup>+6.88</sup> <sub>-6.88</sub>		4.07	420.11 <sup>+7.12</sup> <sub>-7.17</sub>	1.33 <sup>+0.008</sup> <sub>-0.008</sub>	9.47 <sup>+8.34</sup> <sub>-8.34</sub>		1.12	

# Results

Goodness-of-fit comparison table

Galaxy	NFW	Burkert	Kroupa		NFW	Burkert	diet-Salpeter	
			Rindler ( $n = 1$ )	Rindler ( $n \neq 1$ )			Rindler ( $n = 1$ )	Rindler ( $n \neq 1$ )
D 154	1.06	0.43	2.10	1.77	1.09	0.39	2.03	1.75
I 2574	2.40	0.69	4.74	0.98	5.49	1.16	5.53	1.32
N 2366	3.00	1.34	3.48	3.27	2.94	1.34	3.29	3.23
N 2403	0.82	1.60	11.80	2.27	1.08	1.37	9.29	2.55
N 2841	1.29	3.02	76.10	0.90	0.49	0.93	44.30	1.23
N 2903	3.43	2.01	71.80	7.99	3.73	1.23	74.10	8.34
N 2976	6.31	1.00	8.05	1.15	10.49	5.11	10.70	10.60
N 3031	4.96	5.12	9.54	5.12	22.76	4.37	22.80	20.30
N 3198	4.75	2.49	14.00	7.77	7.86	0.93	11.50	9.54
N 3521	5.66	5.21	6.14	5.82	12.48	11.5	7.54	6.84
N 3621	1.45	5.59	11.20	0.87	1.44	3.64	3.29	1.24
N 4736	1.34	1.30	20.50	5.76	1.30	1.29	8.74	4.11
N 5055	5.09	4.11	5.09	4.96	15.74	3.47	15.80	13.50
N 6946	1.21	1.14	1.29	1.25	4.42	2.49	4.42	2.47
N 7331	8.63	7.06	8.65	8.21	29.60	15.6	29.60	26.00
N 7793	4.07	7.88	10.50	4.48	5.01	4.30	5.52	5.31
N 925	3.68	1.19	3.74	1.35	5.64	2.54	5.72	2.30

# Results

## Summarizing:

- Fits for  $n = 1$  are poor, since most of the  $\chi_{\text{red}}^2$  are bigger than unity. However, the most important problem is the Rindler acceleration parameter does not converge to a single value. The computed parameter is in the interval  $0.93_{-0.44}^{+0.01} < a < 9.57_{-0.06}^{+0.06}$  in Kroupa's model, to account for a difference of an order of magnitude, whereas  $0.62_{-0.10}^{+0.10} < a < 7.79_{-0.06}^{+0.06}$  in diet-Salpeter's model, in a similar fashion as in the previous model.
- When the power-law parameter  $n$  is set free, the fits become better, and by comparing them Kroupa's did much better in 14 (out of 17) cases than diet-Salpeter's. However, the goodness-of-fit test does not render acceptable results since some galaxies present very high  $\chi_{\text{red}}^2$  values. The parameters spreads are large too: For Kroupa's model:  $77.96_{-1.75}^{+1.74} < a < 32294.30_{-212.01}^{+211.39}$ , accounting for two orders of magnitude in difference and  $0.002 \sim n < 2.14_{-0.044}^{+0.045}$  that yields a three orders of magnitude difference. For diet-Salpeter parameters the spreads are similar.

# Results

## Summarizing:

- When the mass-to-light ratio is set free, the fits are better than previous models when  $n$  is a free parameter, but again the spread in parameters is high:  $10 \sim a < 19022_{-154}^{+153}$ , accounting for a three orders of magnitude difference and  $0.002 \sim n < 1.52_{-0.009}^{+0.009}$ , a difference of two orders of magnitude.
- Comparing these results with those of NFW and Burkert, the  $\chi^2$ -test favors firstly Burkert and then NFW profiles over Rindler's modify gravity. The tendency is clearer for the standard Rindler ( $n = 1$ ) that fits worst than both NFW's and Burkert's profiles for the Kroupa and diet-Salpeter stellar mass models. The generalized Rindler model ( $n \neq 1$ ) for the diet-Salpeter stellar mass model results are slightly better than the Kroupa's model, but still the NFW profile fits better for 9 galaxies (out of 17) and Burkert profile achieves a better fit for 12 galaxies (out of 17) than the generalized Rindler model.

- 1 Modified Gravity at Large Distances
- 2 Rotation Curves
- 3 Results
- 4 Conclusions**

# Conclusions

- We have tested the idea that a modification of the Newtonian potential stemming from a Rindler acceleration that modifies gravity, as

$$\phi_i(r) = -G\frac{M}{r} + a r^n \quad \Rightarrow \quad v_T^2(r) = \Upsilon_\star v_\star^2 + v_G^2 + a|\vec{r}|^n \quad (15)$$

- The overall conclusion is that although the Rindler modified gravity fits are achievable for the considered galaxies, in many cases they show high  $\chi_{\text{red}}^2$  values, and a high spread in the Rindler parameters  $(a, n)$  that points for an inconsistent model. Furthermore, the standard dark matter profiles (NFW and Burkert) or the alternative BDM model do a better job to fittings of the rotation curves.