Primordial fluctuations from deformed quantum algebras

Sanjeev Seahra (with A Day and I Brown)

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Inflation as a probe of quantum gravity

Quantum gravity effects on perturbations

Deformed quantum algebras

QFT in flat space

Primordial fluctuations

Conclusions

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 - the spatial distribution of these progenitors are assumed to fixed by quantum processes during inflation
 - quantum gravity effects should be imprinted on distribution of galaxies in the sky

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hence if quantum gravity effects manifest themselves distances less that some "exotic physics" length scale M_{\star}^{-1} , one expects modifications of perturbative dynamics when

 $\frac{\text{exotic physics length scale}}{\text{phys. wavelength}} \sim g = \frac{k}{M_{\star}a} \gg 1$

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- no calculable quantum gravity theories available \Rightarrow consider phenomenological models that alter:
 - \Box classical dynamics for large g (via semiclassical effective theories, modified dispersion relations, etc.), or
 - □ quantum dynamics for large g (e.g. deformed quantum algebras) arXiv:1311.1059-3/7

Deformed quantum algebras

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assume that Heisenberg algebra of quantum operators gets modified at high energy:

$$[\hat{q}, \hat{p}] = i \quad \mapsto \quad [\hat{q}, \hat{p}] = if(\beta \hat{p})$$

where β is a dimensionful parameter that defines a scale where deformations become important

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- Want to recover ordinary physics for small momenta \hat{p} and have invariance under $(\hat{q}, \hat{p}) \mapsto (-\hat{q}, -\hat{p})$:

$$f(0) = 1$$
 $f(x) = f(-x)$

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arXiv:1208.5761 (Husain, Kothwala, SSS): studied effects of DQAs on non-interacting QFT in flat space

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assumed that the commutator of field amplitude and momenta modified in Fourier space:

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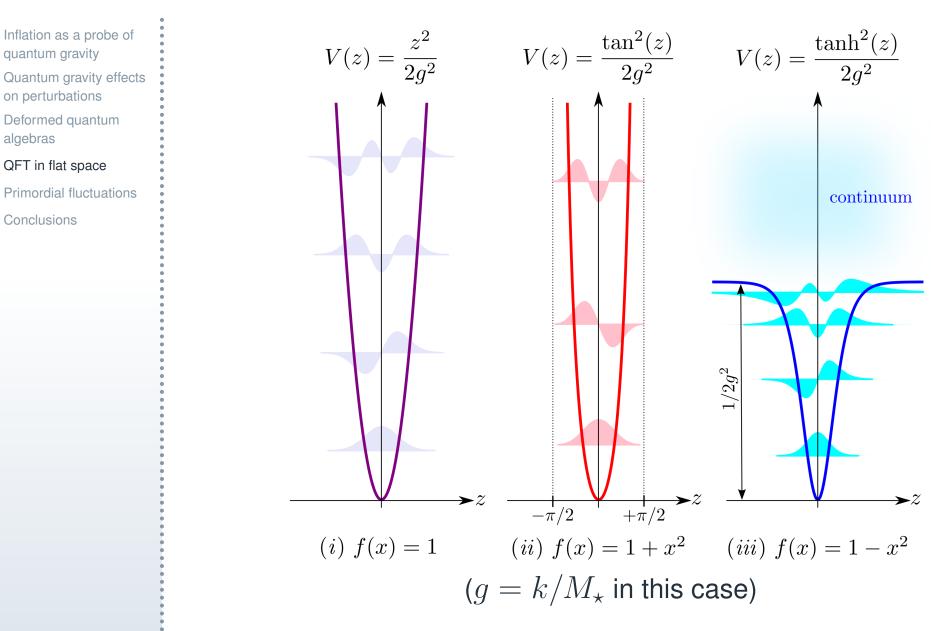
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- number of interesting effects:
 - □ modified propagator
 - □ modified dispersion relations
 - □ field energy no longer quantized in units of $\hbar \omega$, energy levels follow from solution of a 1D Schrödinger bound state problem



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$$\epsilon = \frac{\text{inflationary Hubble parameter}}{\text{exotic physics scale}} = \frac{H}{M_{\star}}$$

 $\Box \quad k_{\star} = \text{comoving wavenumber of a mode with physical wavelength } 2\pi/M_{\star} \text{ at start of inflation}$

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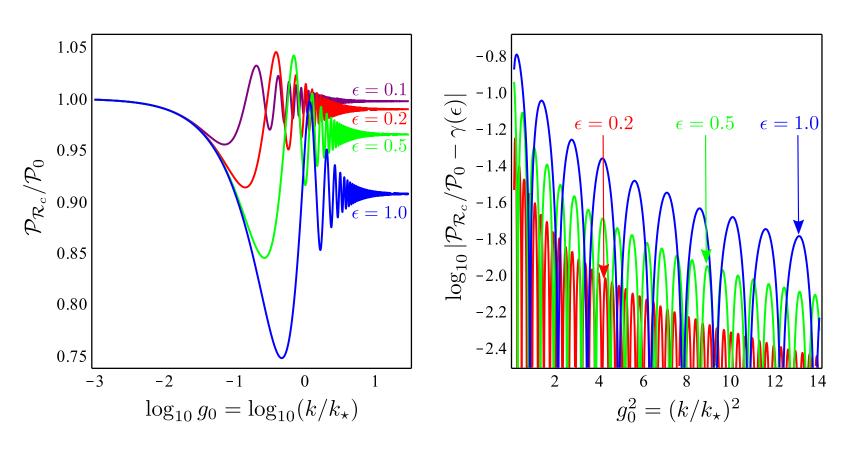
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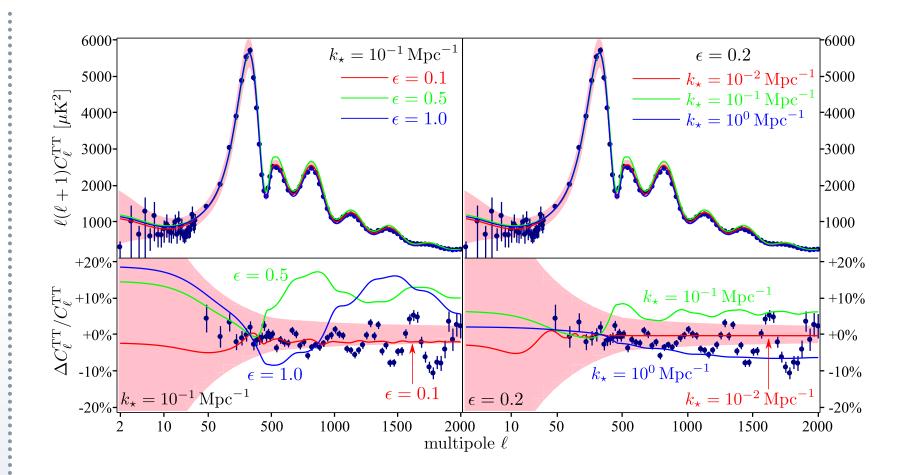
$$k_{\star} \sim \frac{6 \times 10^{-6}}{\mathrm{Mpc}} \epsilon^{-1} \left(\frac{E_{\mathrm{inf}}}{10^{16} \,\mathrm{GeV}}\right) \left(\frac{e^{65}}{e^{N}}\right) \left(\frac{100}{\mathcal{G}}\right)^{1/12}$$

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comoving curvature perturbation results for f(x) = 1 + x²
\$\mathcal{P}_0\$ is ordinary slow-roll result—recovered for small \$k\$
DQA effects \$\propto \epsilon^2 = H^2/M_{\star}^2\$

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I CMB power spectrum results for $f(x) = 1 + x^2$

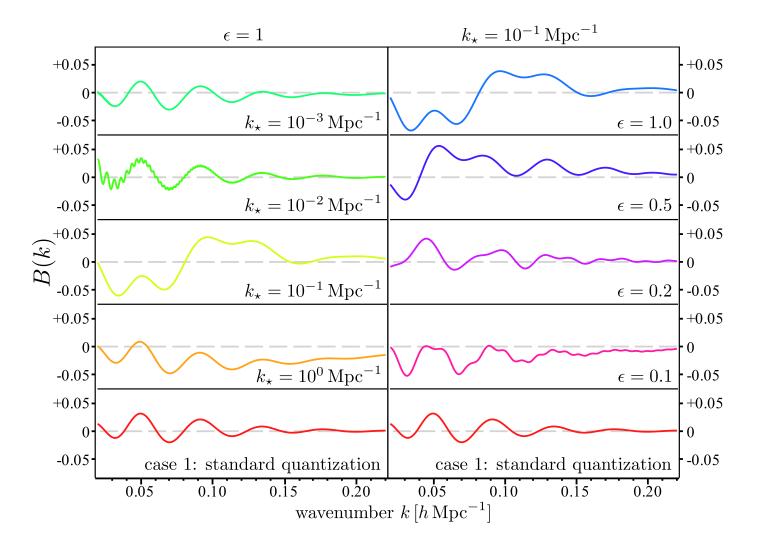
high multipoles can distinguish between various values of $\epsilon = H/M_{\star}$

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Baryon acoustic oscillations (BAOs) for $f(x) = 1 + x^2$

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- effects were $\mathcal{O}(H^2/M_\star^2)$ (for a specific case)
- may be possible to see effects in CMB or BAOs if H/M_{\star} not too small
- **still to do:** parameter estimation/limits on ϵ and k_{\star}